



## The Fermi haze from Dark Matter Annihilation and Anisotropic Diffusion

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arXiv:1102.5095 (accepted by ApJ),

ApJ 717,825,(2010) (arXiv:0910.4583) (G. Dobler, D. Finkbeiner, IC, T. Slatyer, N. Weiner), arXiv:0911.4954 (IC, N. Weiner)

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## Outline

- Fermi haze/Fermi bubbles a true signal. Lower latitudes uncertainties
- Dark Matter case (models that work)
- Anisotropic diffusion of CRs in the Galaxy
- Conclusions



## The first Fermi haze template









One needs to be very careful for small (but significant in the interpretation) caveats with using templates.



Su, Slatyer and Finkbeiner work



50

0

-50

50

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ApJ 724, 1044 (2010) (arXiv:1005.5480)

### Fermi Bubbles (Su, Slatyer, Finkbeiner) (SFD+disk+uniform+bubbles)



(Updated) Fermi haze (G. Dobler, IC, N.Weiner)

We used the Fermi gamma-ray map at 0.5-1 GeV as a background galactic template+uniform + haze(modeled by GALPROP Dark Matter IC signal)



X-shape that could indicate an over-subtraction (of pi0 gamma-rays) in the SFD template scheme. The pi0 to dust column ratio needs not to be constant (for instance a source heating up the region and giving also the signal in X-rays). So while at the high latitudes the signal is clear (even some edge effect at high latitudes is confirmed) BUT at lower latitudes the template selection is important.

## What about Dark Matter?

The DM smooth halo has an approximately Spherical distribution, a possible candidate.

DM can explain the haze signal (WMAP + Fermi) as has been shown in arXiv: 0911.4954 (IC + N. Weiner) based on solely energetic/spectral arguments (XDM electrons with local annihilation BF ~ 100 (~50 at the haze region)).



Leptophilic DM models can explain the signal. Models that annihilate to taus or have large BRs to hadrons can not explain the angular morphology of the signal.

## Anisotropic and inhomogeneous CR diffusion in the ISM

Propagation equation:

$$\frac{\partial \psi}{\partial t} = \frac{\partial (b\psi)}{\partial E} + \overrightarrow{\nabla} (D\overrightarrow{\nabla}\psi) + Q \tag{I}$$

 $\psi$  is the CR number density at time  $t\,$  and position  $\vec{x}$ 

*b* : energy loss coefficient (above 5GeV dominated by IC and synchrotron emission).

- D : diffusion constant
- Q : source term

Assuming cylindrical symmetry:

$$\overrightarrow{\nabla}(D\overrightarrow{\nabla}\psi) = \frac{1}{r}\frac{\partial}{\partial r}(rD\frac{\partial\psi}{\partial r}) + \frac{\partial}{\partial z}(D\frac{\partial\psi}{\partial z})$$
(II)

Anisotropic diffusion:

$$\vec{\nabla} (D \vec{\nabla} \psi) = \frac{1}{r} \frac{\partial}{\partial r} (r D_{rr} \frac{\partial \psi}{\partial r} + r D_{rz} \frac{\partial \psi}{\partial z}) + \frac{\partial}{\partial z} (D_{zz} \frac{\partial \psi}{\partial z} + D_{zr} \frac{\partial \psi}{\partial r})$$
(III)

What we will assume is a strong magnetic field perpendicular to the galactic plane in the inner part of the Galaxy.

Random(irreg.) B-field component:  

$$B_{irreg} = B_0 e^{(R_{\odot} - r)/r_1 - |z|/z_1}$$
  
 $R_{\odot} = 8.5 kpc$ 

Ordered B-field component:

$$B_{\rm ord} = B_1 e^{-r/r_2 - |z|/z_2} \times \left(1 + K e^{-r/r_3 - |z|/z_3}\right)$$

What remains is to relate the elements of the diffusion tensor to the magnetic field.

$$D \propto \lambda_{sc} \propto r_{gyr} \propto B^{-1}$$

Also assuming that the ordered field is along z-axis and much stronger than the turbulent field we expect:

 $\lambda_{sc_z} \gg \lambda_{sc_r}$ 

Following formulation developed by Parker (1965)

- v : frequency by which CRs scatter off from their spiral orbit
- $\Omega \gg v :$  in the central part of the Galaxy
- $\Omega \ll v \,$  : far from the galactic center

we have:

$$D_{zz} \propto B_{tot}^{-1} \left( \frac{v^2 + \frac{q^2 B_z^2}{\gamma^2 m^2 c^2}}{v^2 + \frac{q^2 B_{tot}^2}{\gamma^2 m^2 c^2}} \right)$$

setting:

$$A = \frac{q}{\gamma m c \mathbf{v}}$$

we get:

$$D_{zz} \propto B_{tot}^{-1} \left(\frac{1 + A^2 B_z^2}{1 + A^2 B_{tot}^2}\right)$$

#### (extreme example)



$$\frac{D_{rr}}{D_{zz}} = \frac{1 + A^2 B_r^2}{1 + A^2 B_z^2}, \ \frac{D_{rz}}{D_{zz}} = \frac{D_{zr}}{D_{zz}} = \frac{A^2 B_r B_z}{1 + A^2 B_z^2}$$

### Thus one can get:



So with annihilating DM and specific assumptions on anisotropic and inhomogeneous diffusion we CAN fit the Fermi haze morphology spectrum and amplitude.

Different assumptions for the B-field can have apart from different synchrotron maps, different IC maps.

Model	$B_{\rm ord}$ Formula	$\begin{array}{c} B_0 \\ (\mu \mathrm{G}) \end{array}$	$r_1$ (kpc)	$\begin{array}{c} z_1 \ ( ext{kpc}) \end{array}$	$\begin{array}{c} B_1 \\ (\mu \mathrm{G}) \end{array}$	K	$r_2$ (kpc)	$\begin{array}{c} z_2 \ (\mathrm{kpc}) \end{array}$	$r_3$ (kpc)	$z_3$ (kpc)
1	$B_1 e^{-r/r_2 -  z /z_2} \times (1 + K e^{-r/r_3 -  z /z_3})$	3	7	4	8	10	7	2	0.8	10
2	$B_1 e^{-r/r_2 -  z /z_2} \times \left( 1 + K e^{-(r/r_3)^2} \sqrt{\cos( z /z_3 \times \pi/2)} \right)$	3	5	4	10	11	5	4	1	40
3	$B_1 e^{-r/r_2 -  z /z_2} \times \left(1 + K e^{-(r/r_3)^{1.5} -  z /z_3}\right)$	3	10	2	10	6	10	3	1.2	20
4	$B_1 e^{-r/r_2 -  z /z_2} \times \left(1 + K e^{-(r/r_3)^{1.5} - ( z /z_3)^{1.5}}\right)$	3.7	5	2	12.5	8	7	5	2.5	20
5	$B_1 e^{-r/r_2 -  z /z_2} \times (1 + K e^{-r/r_3 -  z /z_3})$	3.7	5	2	3.7	12	5	2	2	6

# Anisotropic diffusion of CRs can have a strong effect on the IC map from annihilating DM:







Apart form getting the right morphology for the Fermi (and WMAP) haze, and having good spectral agreement, we also have agreement with local CRs and background gamma-rays.



DM Contribution

## Conclusions

- The Fermi haze is a signal of a population of harder spectrum electrons (seen before only at microwave) that "conventional" sources of electrons such as middle-aged pulsars can not explain
- DM with Anisotropic and inhomogeneous diffusion may be the answer
- Astrophysical explanations including MSPs (Malyshev, Cholis, Gelfand ApJ 722, p.1939-1945 (2010)), or a strong AGN activity may be in order(Guo & Mathews arXiv:1103.0055), strong Galactic wind (Crocker&Aharonian PRL 106:101102,2011), 2nd order Fermi acc. (Mertsch&Sarkar arXiv:1104.3585 (PRL))
- Need further modeling and calculations on the signal in order to better understand the gamma-ray backgrounds AND work out the signals from the possible sources
- neutrinos can be drastically different among the different models

Thank you

## Additional slides

## Magnetic field profiles



Yusef-Zadeh & Morris (1987), Morris & Yusef-Zadeh (1989), Morris (2007), have suggested mag. fields up to few mG in large non-thermal radio filaments (with widths of pc and lengths ~ 50pc). Beck (2008) suggested 0.5 mG. Those non-thermal filaments seen by VLA are directed perpendicular to the disk plane, and are probes of the general B-field properties, suggesting a predominantly bipolar field extending ~200pc in r (Nord et. al. (2004)).

Also arguments of CR cooling by synchrotron radiation in the inner 500pc have been used to avoid over-production of gamma-rays by ICS.





Residual and the fit from the Anisotropic Galprop model

### The gamma-ray sky







### Anisotropic diffusion assumptions tests



408 MHz

### Millisecond pulsars & DM



DM annihilating to  $W^+W^-$  with a thermal relic cross-section.

Need 3x10^4 MSPs in the galactic halo! (significant implications about the evolution of the Milky way)

ApJ 722, p.1939-1945 (2010) (arXiv:1002.0587)



Probe a distribution of hard-spectrum electrons, (steady state diff. spectrum of  $\frac{dN_e}{dE}\sim E^{-2}$  )

Fermi haze: inverse Compton scattering WMAP haze: synchrotron radiation

Non-trivial morphology of the Fermi haze (template:bivariate Gaussian) The source(s) responsible for the signal must explain both spectra AND the non-disk-like morphology

### Young pulsars, are probed still pretty well by the SNe distribution.



So clearly conventional astrophysical sources with disk-like distributions, CAN NOT explain the Fermi haze signal.