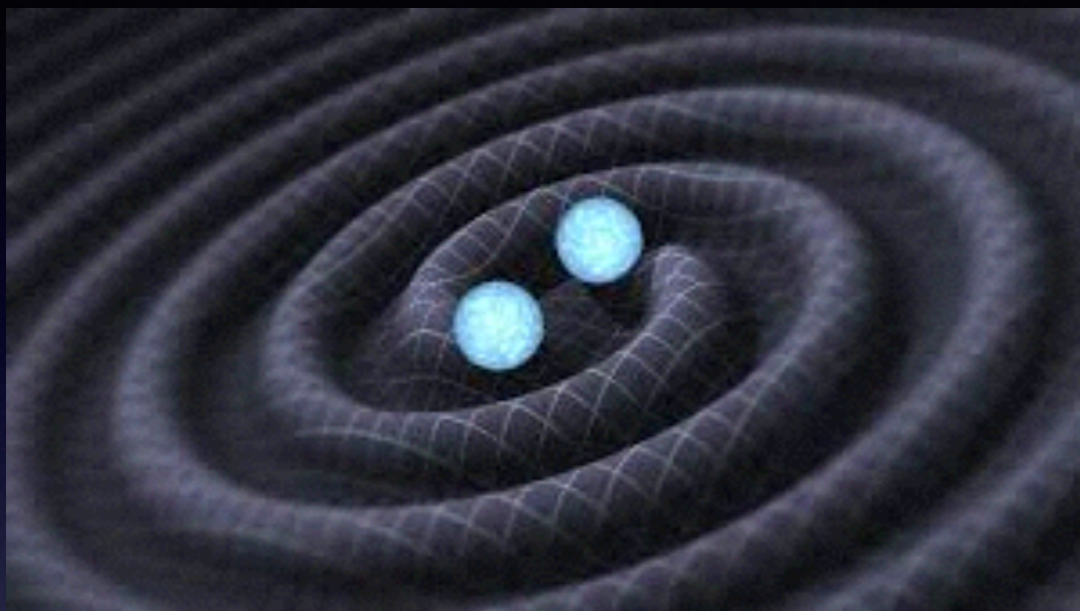


Have we detected Dark Matter with LIGO?



Livingston LA



Hanford WA

aLIGO: PRL, 116.061102; arXiv:1602.03840; arXiv:1602.03847;
S. Bird, I.C, J. Munoz, Y. Ali-Haimoud, M. Kamionkowski, E. Kovetz, A. Raccanelli and A. Riess (JHU) PRL 116.201031,
(arXiv:1603.00464)

I.C., E. Kovetz, Y. Ali-Haimoud, S. Bird, M. Kamionkowski, J.

Munoz, A. Raccanelli PRD 94 084013 (arXiv:1606.07437)

A. Raccanelli, E. Kovetz, S. Bird, I.C. J Munoz PRD 94 023516 (arXiv:1605:01405)

V. Mandic, S. Bird, I.C. (PRL accept.) arXiv:1608.06699

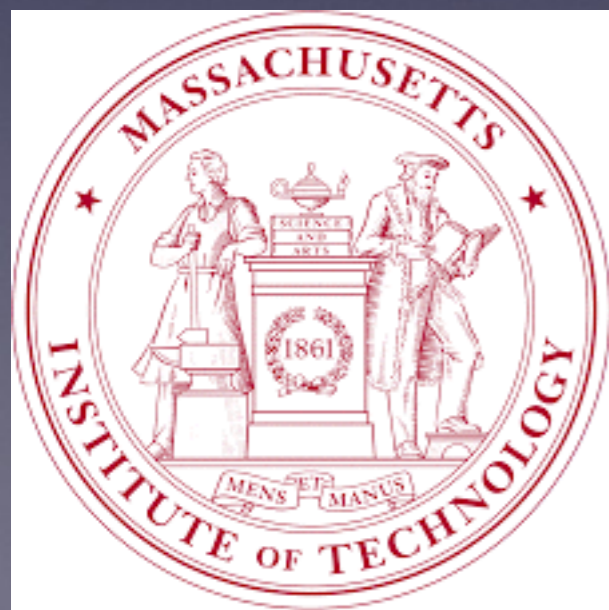
I.C. arXiv:1609.03565

E. Kovetz, I.C., P. Breysse, M. Kamionkowski

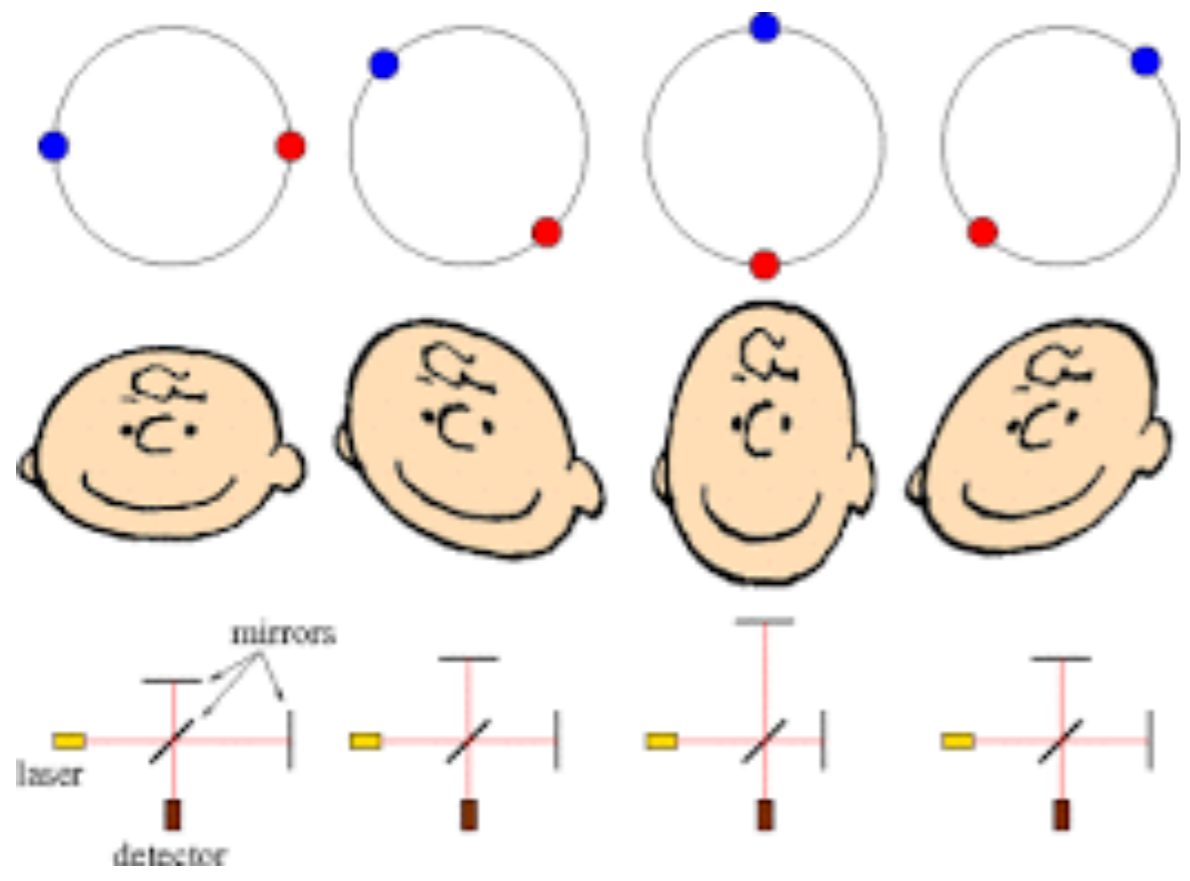
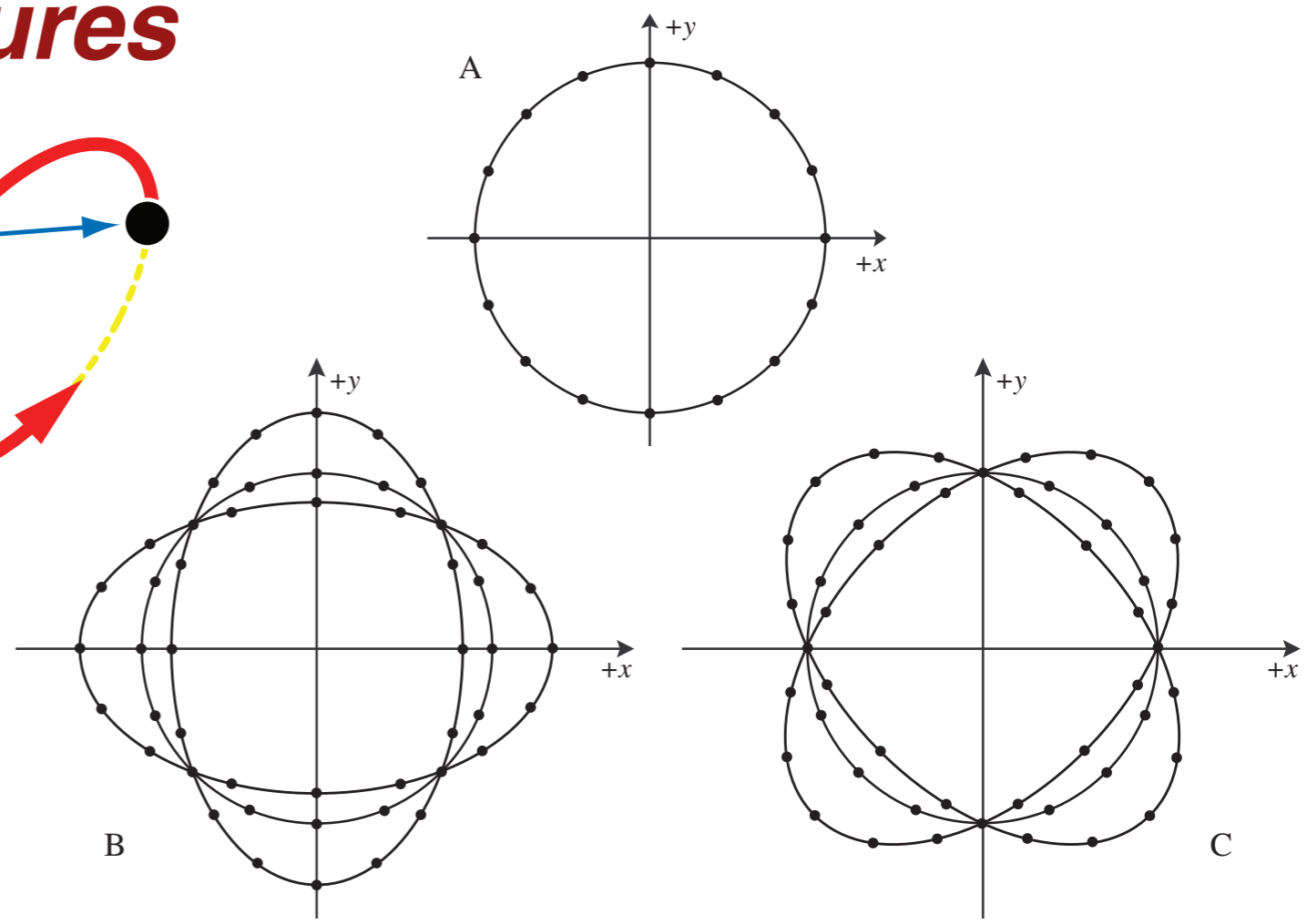
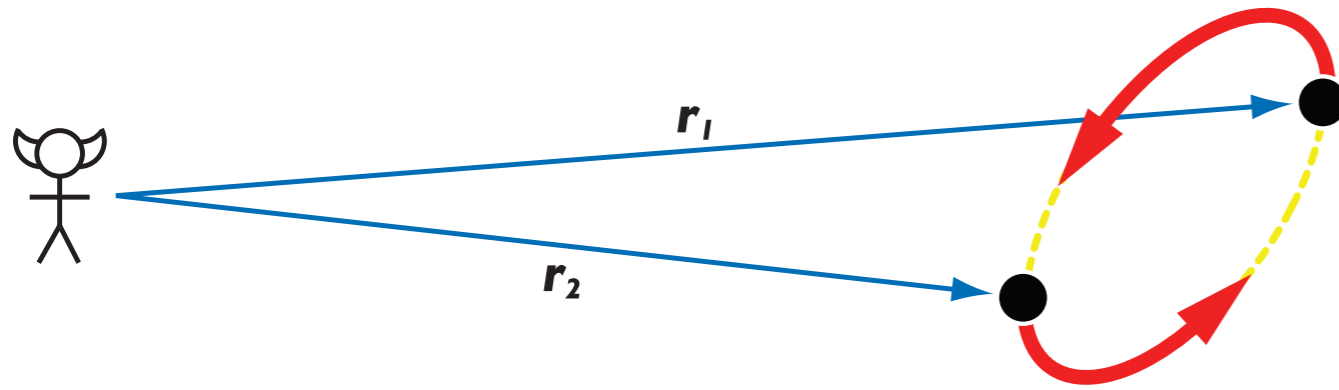
arXiv:1611:01157

Nuclear and Particle Theory Seminar

Ilias Cholis 11/07/2016



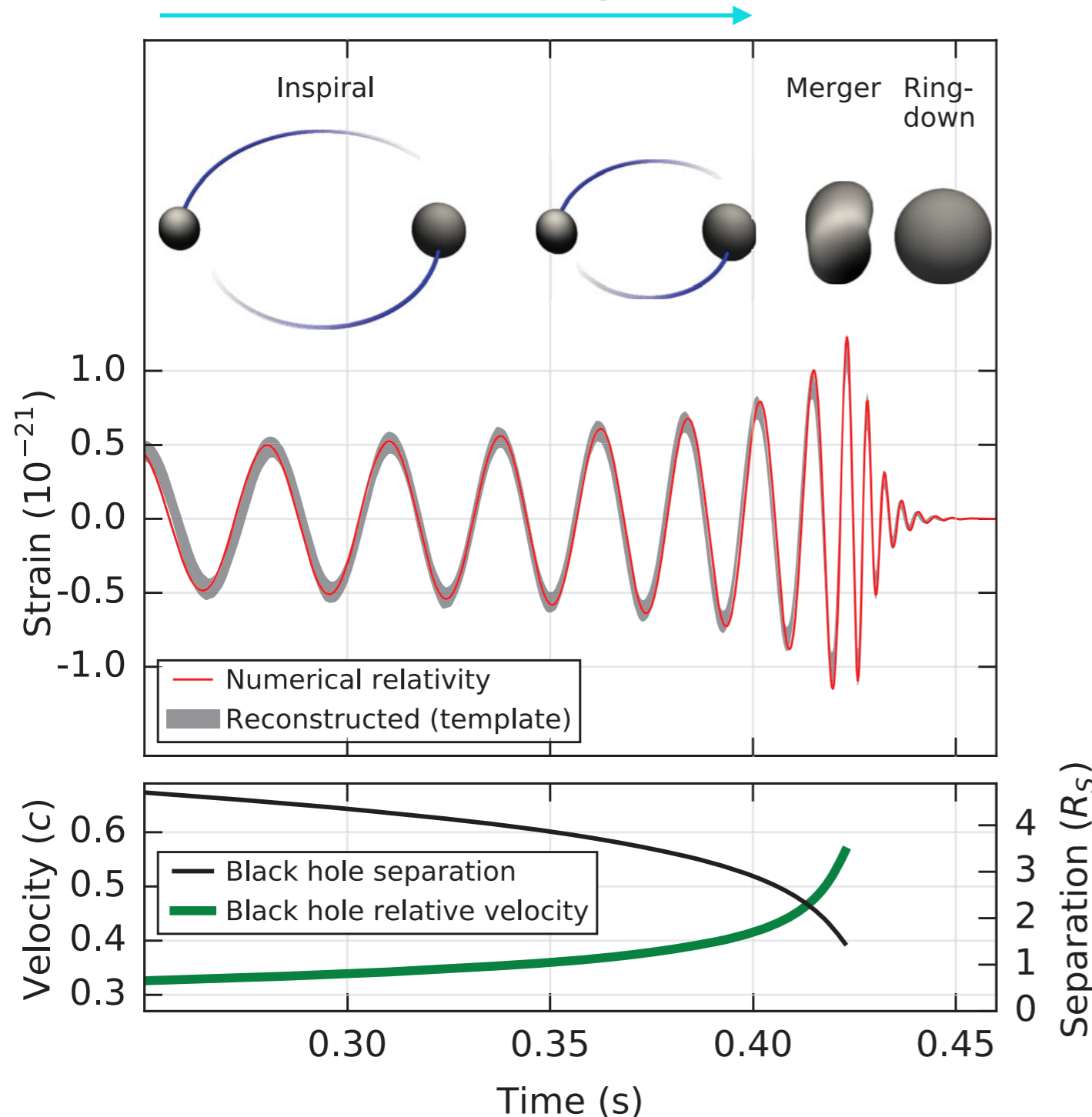
Hand-wavy Pictures



: effect at interferometers

What we expect to observe when two Black Holes coalesce

Frequency Increases with time (“chirp”)



Last Stable Orbit at ~ 3 times the Schwarzschild radius

$$R_{Sch} = \frac{2GM}{c^2}$$

$$R_{Sch}(1M_{\odot}) = 2.95km$$

$$R_{Sch}(36M_{\odot}) = 106km$$

Basic Estimates

GWs travel at the speed of light: $\lambda = c/f$

Take a binary of two compact objects (Kepler's third law):

$$f = \sqrt{\frac{G}{4\pi} \frac{M_{tot}}{a^3}}$$

take $M_{tot} = 20M_{\odot}$ and $a = 500km$ thus $2f \sim 80Hz$

or $\lambda \sim 5 \times 10^3 km$

Earth Size (ground-based Observatories)

take $M_{tot} = 10^6 M_{\odot}$ and $a = 5 \times 10^6 km$ thus $2f \sim 10^{-2} Hz$

or $\lambda \sim 3 \times 10^7 km$

(space-based Observatories)

Basic Scalings

Chirp mass:

Ignoring redshift.

$$M_c = \frac{(m_1 m_2)^{3/5}}{(m_1 + m_2)^{1/5}}$$

Amplitude of signal *during* Inspiral:

$$h_c \sim \frac{G}{c^3} \frac{M_c}{d_L} \left(\frac{G}{c^3} \pi f M_c \right)^{2/3} \rightarrow \text{(for a given freq.): } h_c \sim M_c^{5/3} / d_L$$

(LIGO has a certain freq. range)

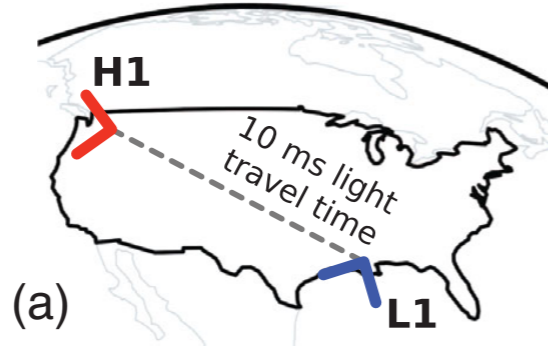
Chirp:

$$\dot{f} = \frac{96}{5} \frac{c^3}{G} \frac{f}{M_c} \left(\frac{G}{c^3} \pi f M_c \right)^{8/3} \rightarrow \text{(for a given freq.): } \dot{f} \sim M_c^{5/3}$$

$$M_c = \frac{c^3}{G} \left(\frac{5}{96} \pi^{-8/3} f^{-11/3} \dot{f} \right)^{3/5} \leftarrow \text{measuring } f \text{ and } \dot{f} \text{ we get } M_c$$

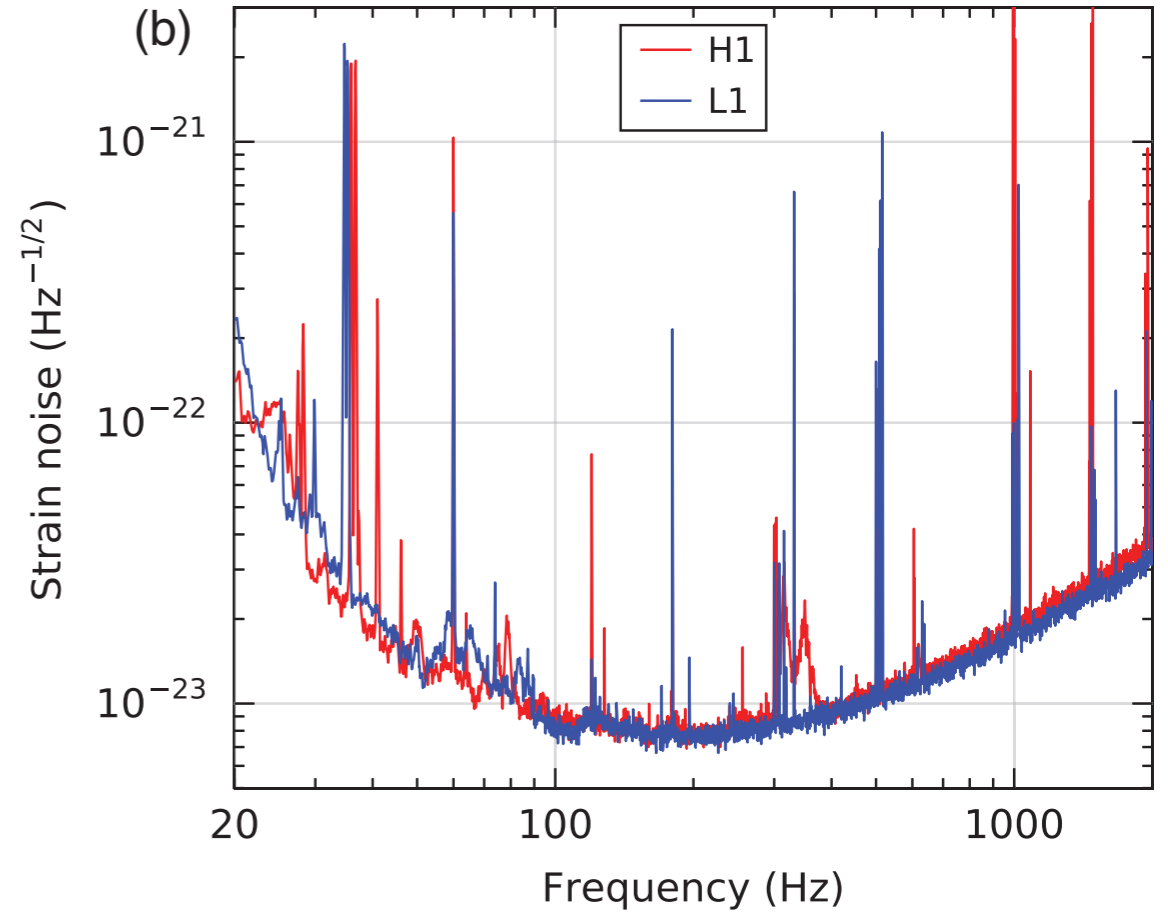
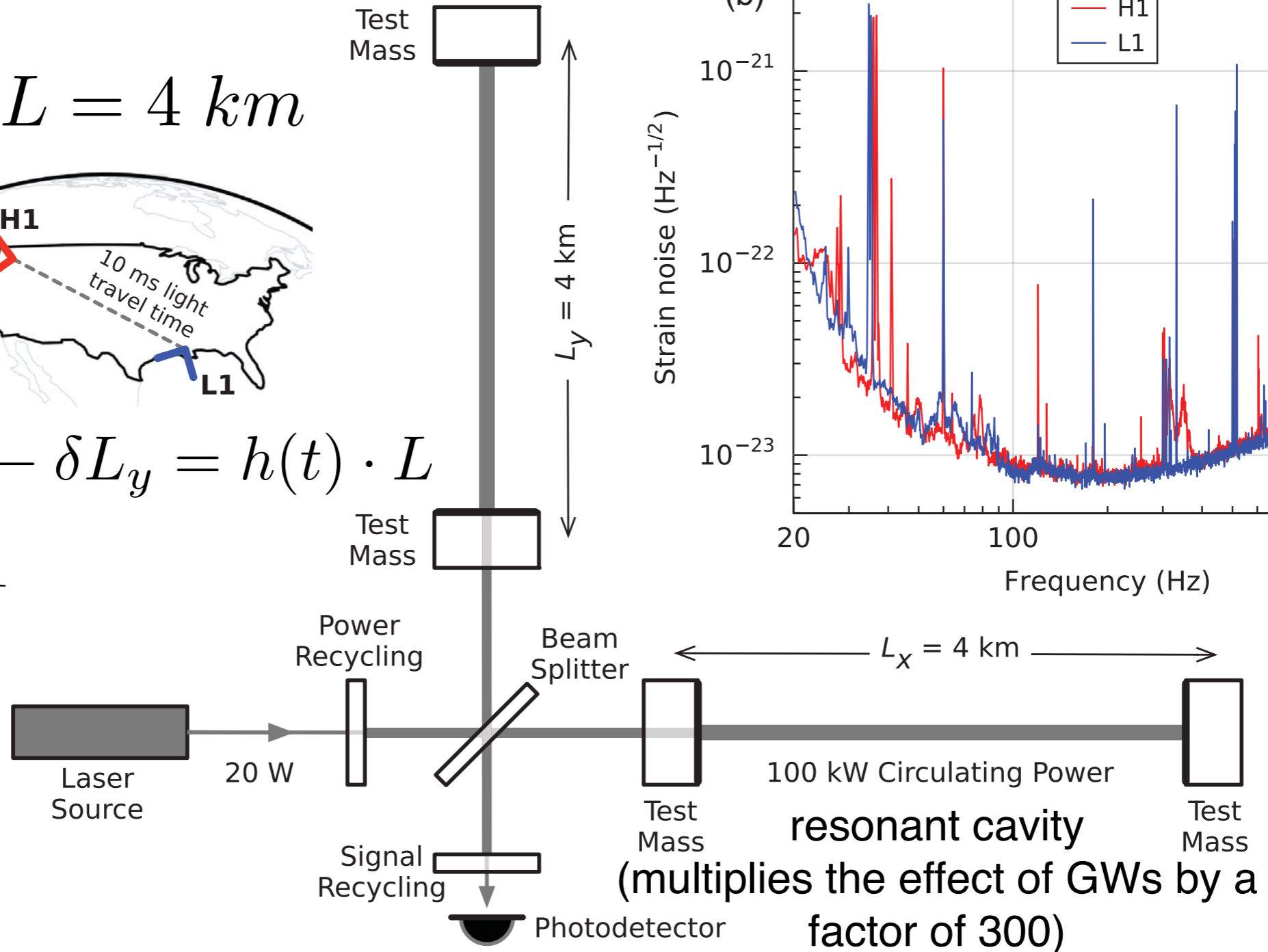
LIGO Detectors

$$L_x = L_y = L = 4 \text{ km}$$



$$\Delta L(t) = \delta L_x - \delta L_y = h(t) \cdot L$$

$$h(t) \simeq 10^{-21}$$



$$h_+(t) = A_{GW}(t)(1 + \cos^2 i) \cos \phi_{GW}(t)$$

$$h_\times(t) = -2A_{GW}(t) \cos(i) \sin \phi_{GW}(t)$$

$$h_k(t) = F_k^+ h_+(t) + F_k^\times h_\times(t)$$

the two grav. wave polarizations

The first ever Gravitational Waves signal detection

On Sept. 14th at 9:50:45 UTC (Coordinated Universal Time), the two detectors of aLIGO observed a gravitational wave signal from the coalescence of two Black Holes. It was observed between 35 and 250 Hz.

The observed Properties are (90 % credible intervals):

$$\begin{aligned} m_1 &= 36_{-4}^{+5} M_{\odot} \\ m_2 &= 29_{-4}^{+4} M_{\odot} \end{aligned} \longrightarrow m_{loss} = 3.0_{-0.5}^{+0.5} M_{\odot} \longrightarrow \dot{E}_{max} = 3.6 \times 10^{56} \text{ erg/s} \quad (\text{instantaneous})$$

$$m_{final} = 62_{-4}^{+4} M_{\odot}$$

$$\alpha = 0.67_{-0.04}^{+0.05}$$

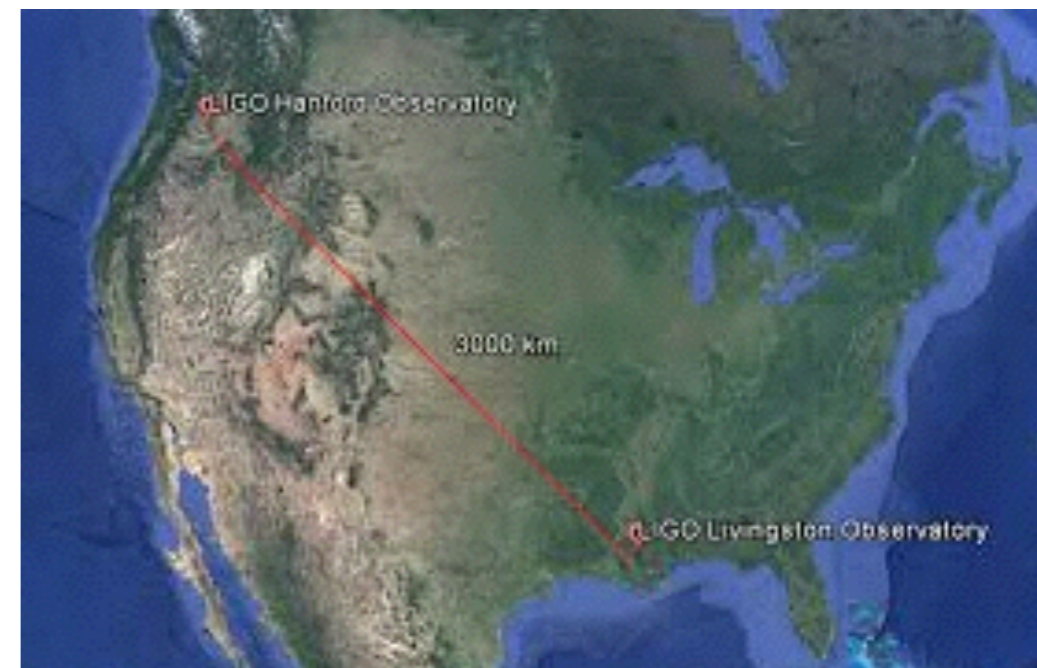
$$\text{final spin: } \alpha = \frac{c |\vec{S}|}{GM^2}$$

$$d_L = 410_{-180}^{+160} \text{ Mpc}$$

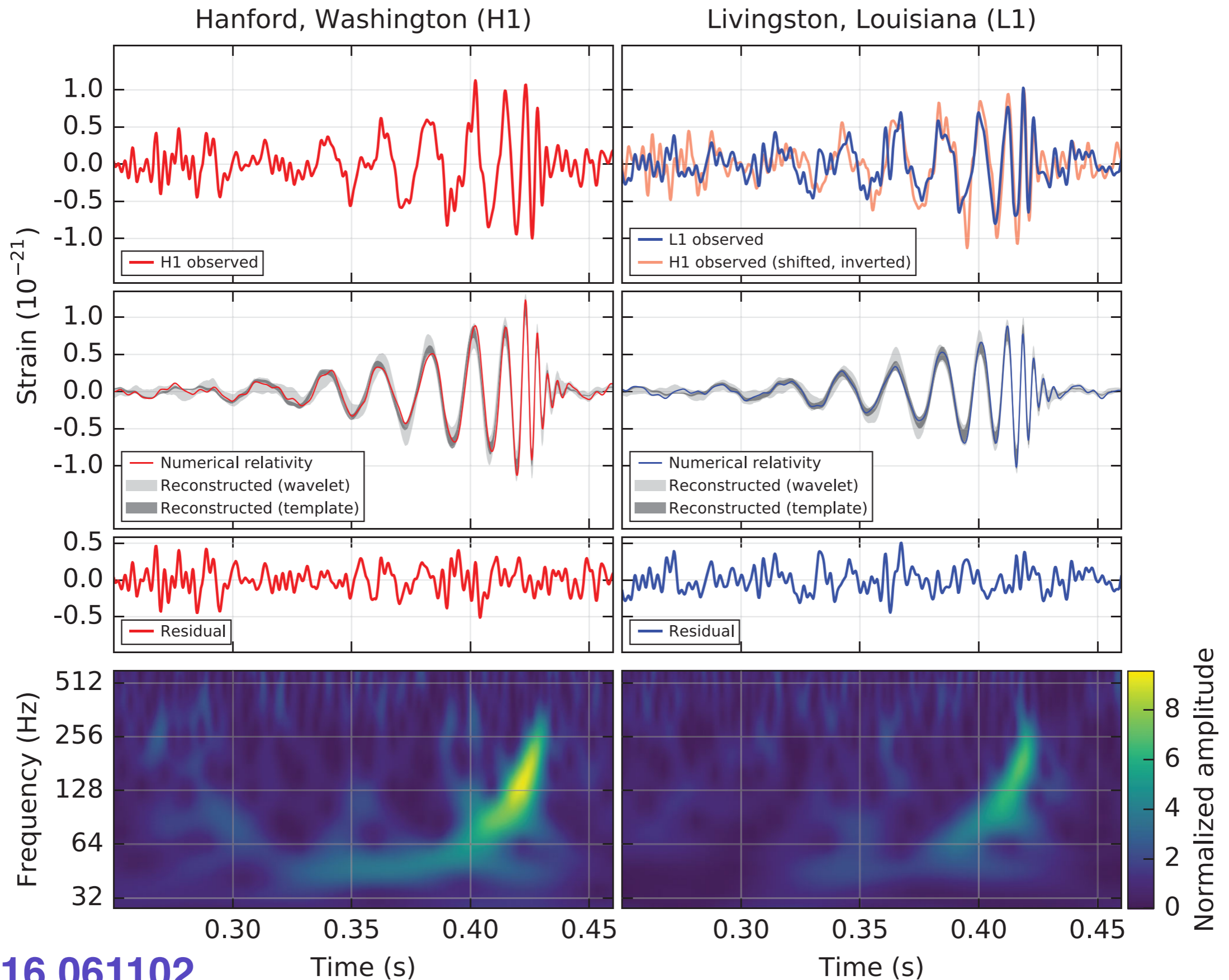
$$z_s = 0.09_{-0.04}^{+0.03} \quad (\text{Planck Cosm. Param.})$$

The event was observed with a time delay of $t_d = 6.9_{-0.4}^{+0.5} \text{ ms}$ between Livingston LA and Hanford WA.

Combined Signal to Noise : 24



The GW150914 event

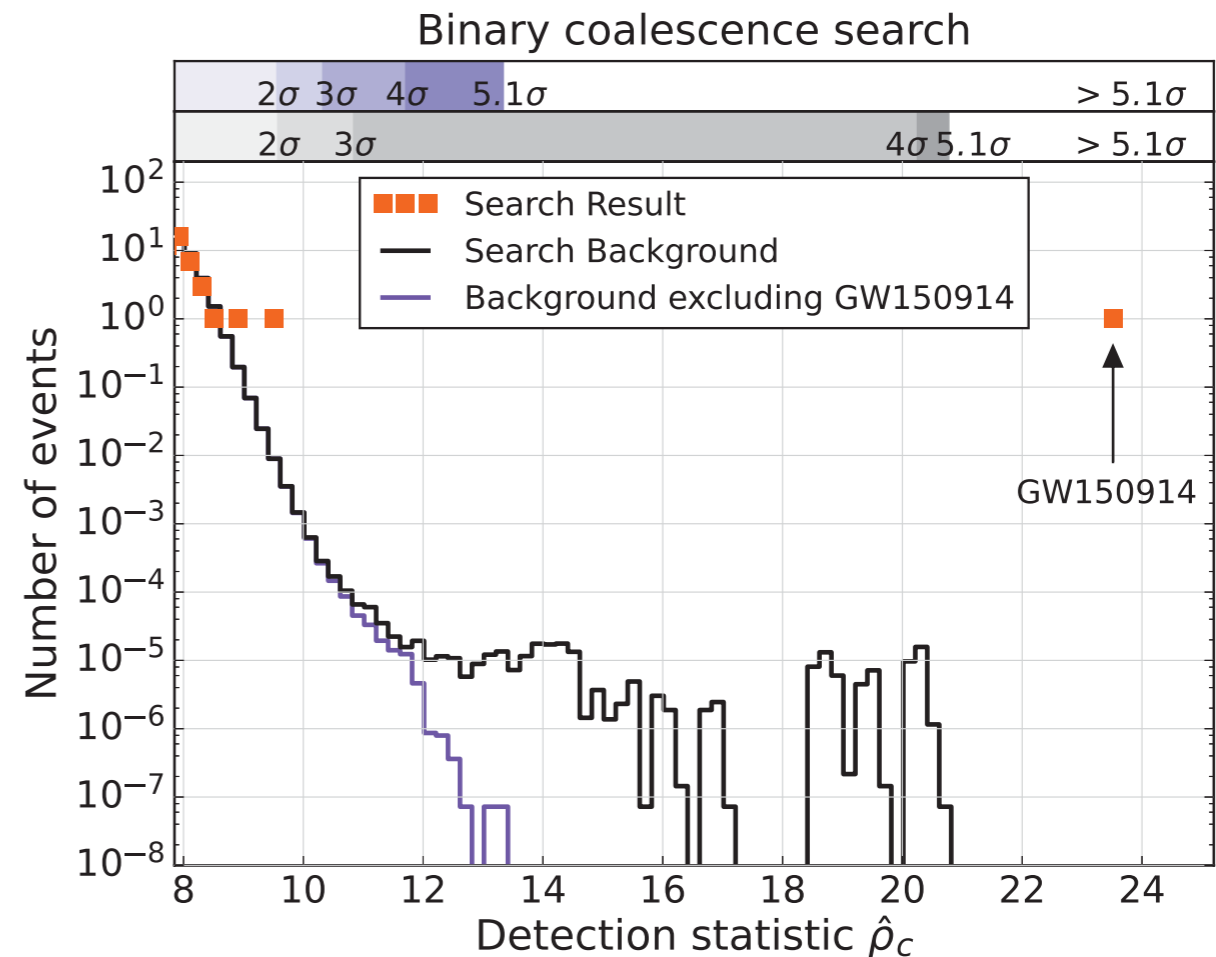
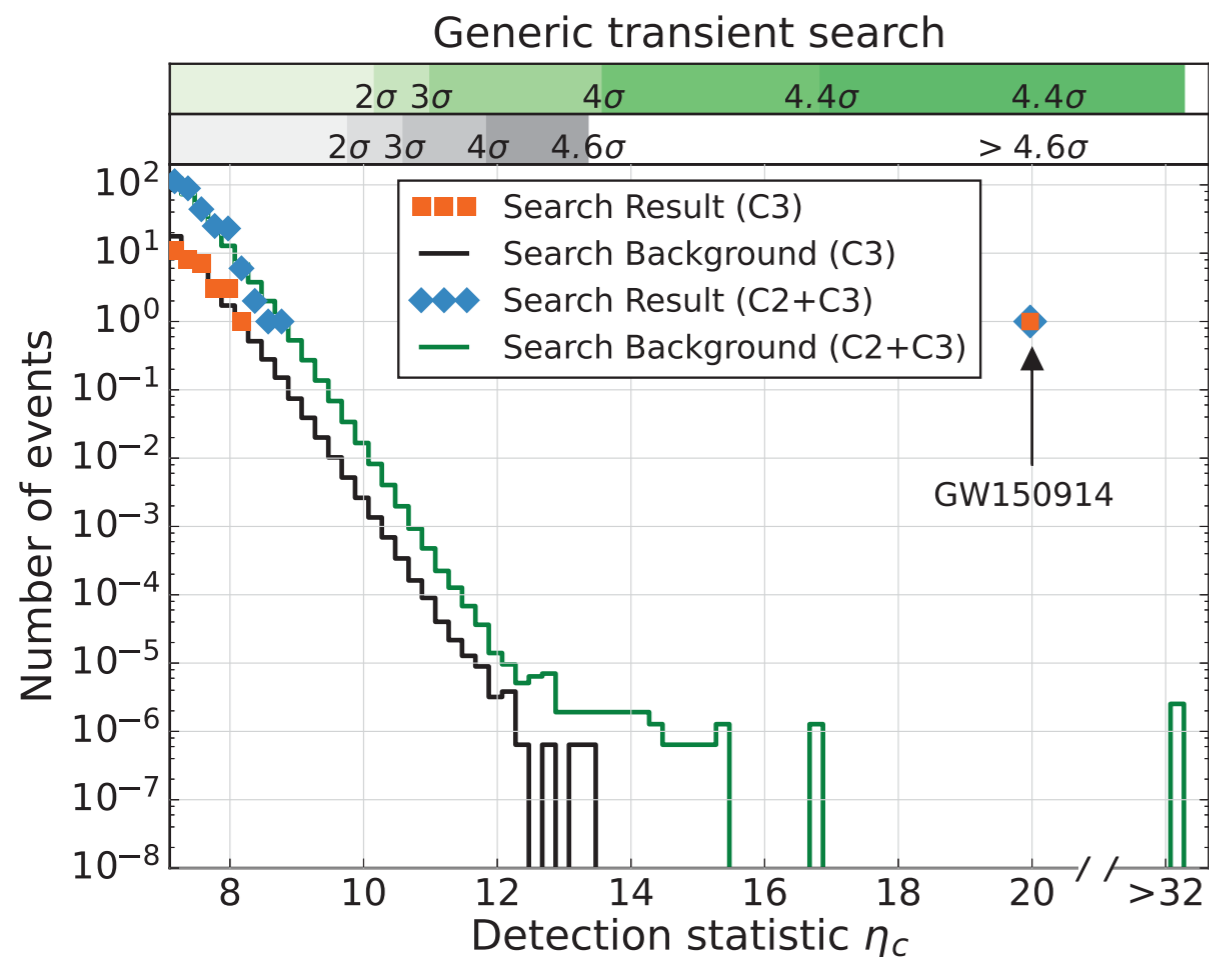


Significance of the event

To understand the significance of a potential GW event LIGO coll. has to estimate the rate at which the detector noise produces a “detection statistic” (likelihood of being a GW) at least as high as the event.

For the generic transient search the false alarm rate is 1 event in 22500 yr $\rightarrow 4.6\sigma$.

For the template binary coalescence events after simulating 203000yr no false alarm $\rightarrow 5.1\sigma$ at least.



Remaining properties

The luminosity distance is correlated to the inclination of the orbital plane to the line of sight θ_{JN} . Total angular momentum \vec{J} .

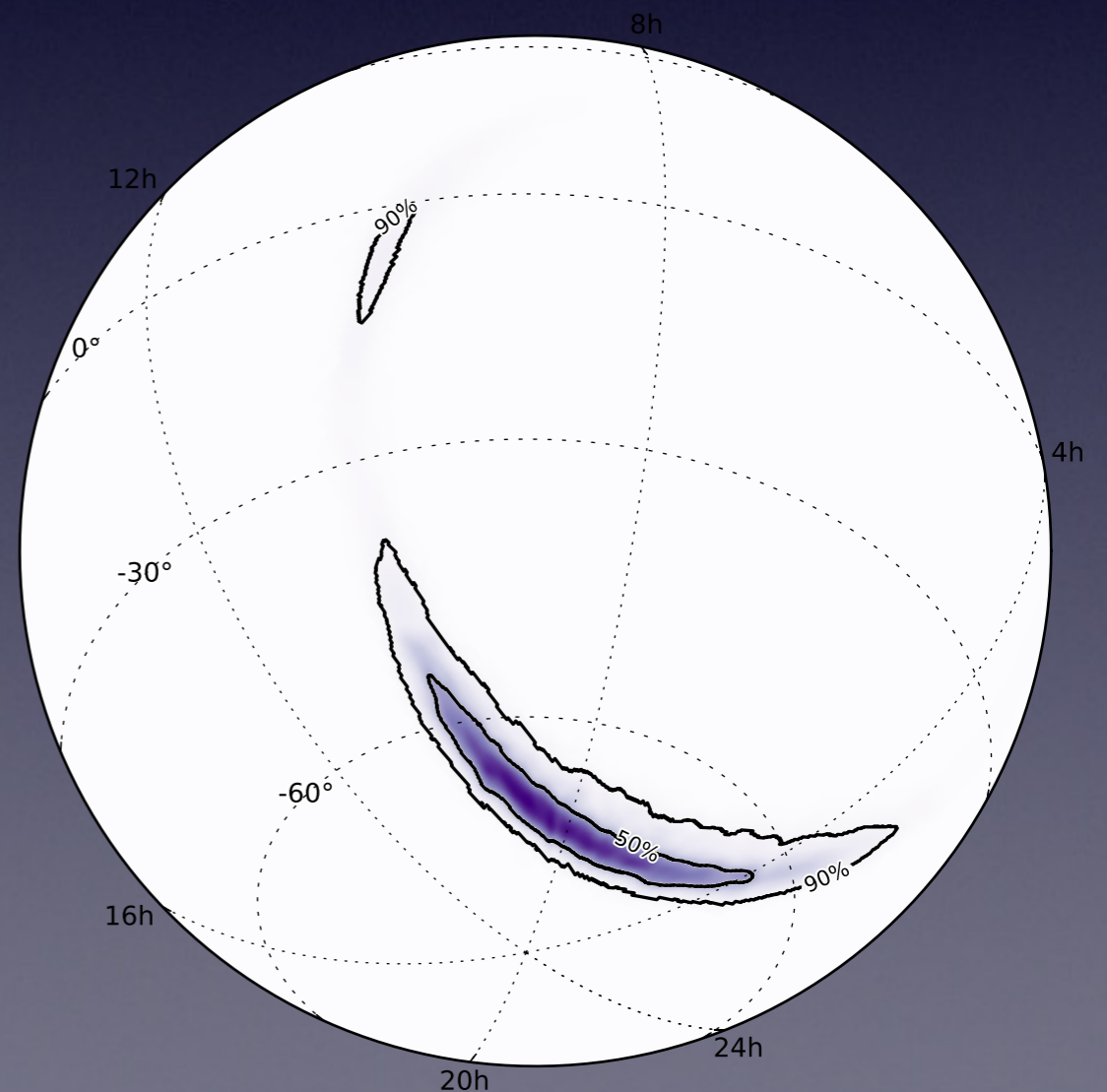
\vec{J} is almost constant during the inspiral.

$45^\circ < \theta_{JN} < 135^\circ$ with a probability of 0.35

50% probability within 140 deg^2

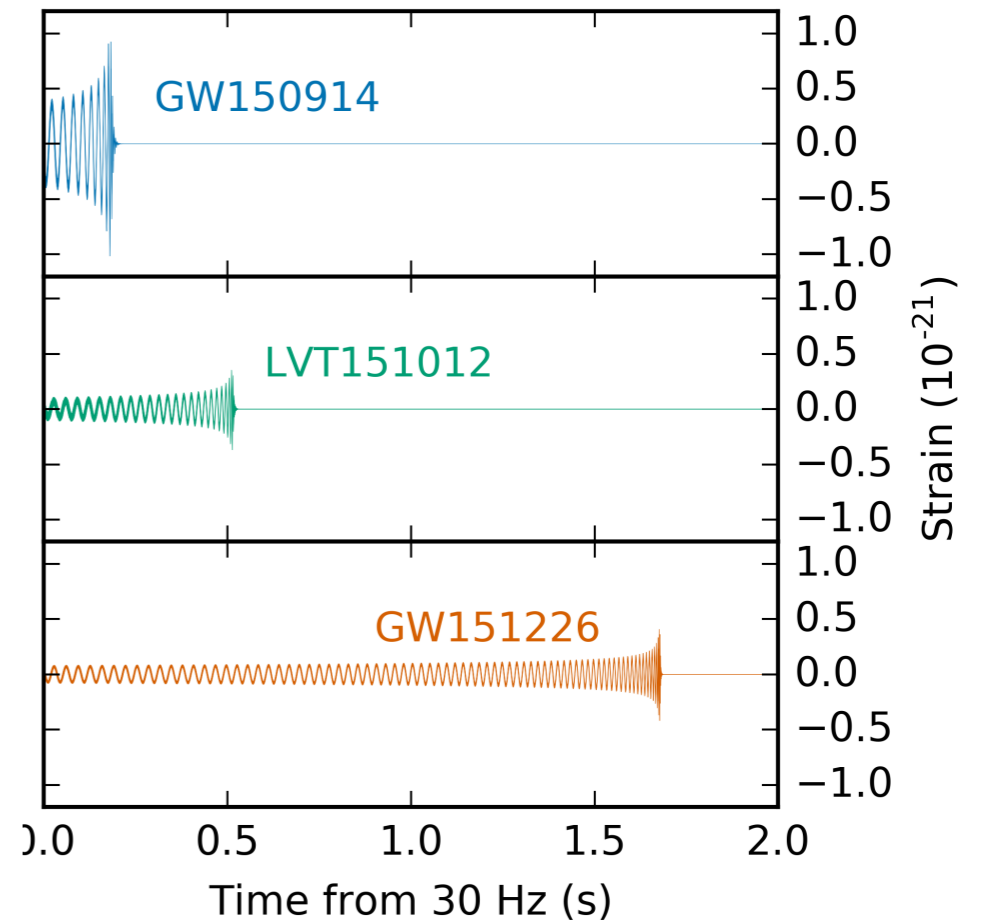
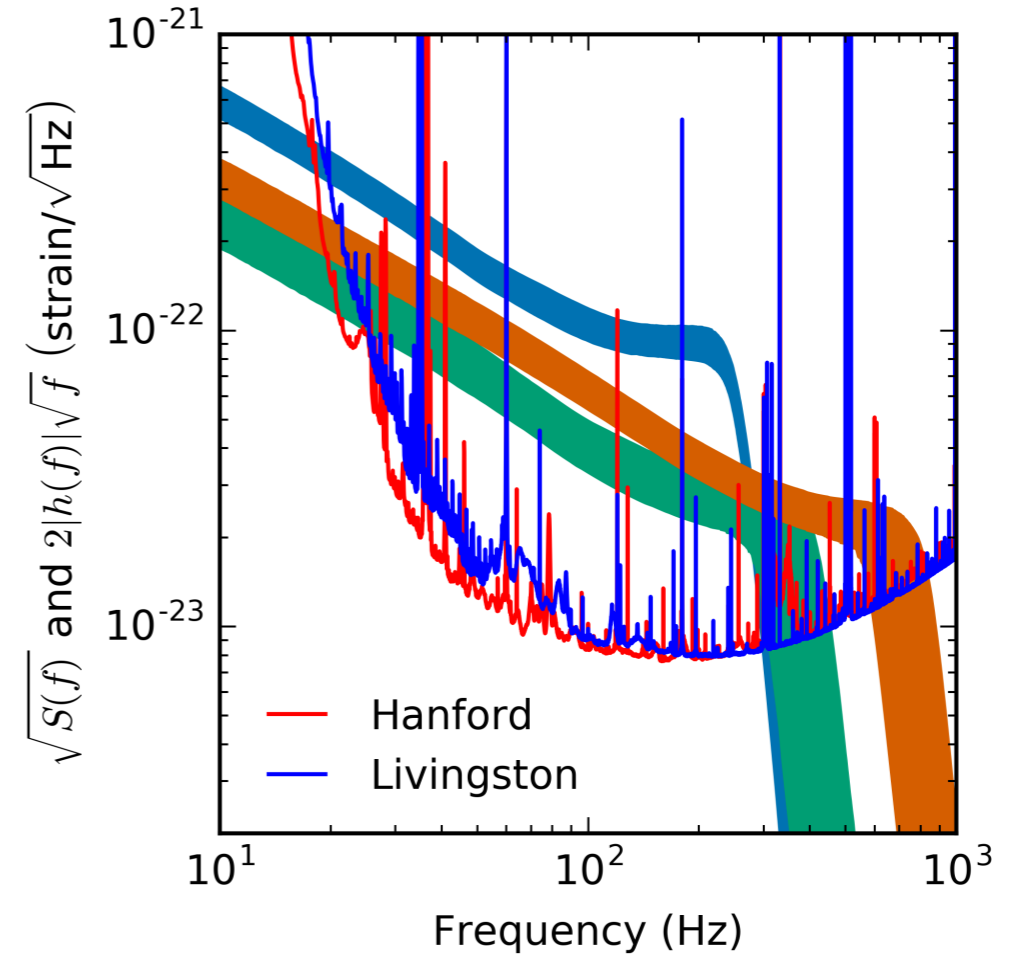
90% probability within 590 deg^2

Searches by EM and neutrino detectors. No evident counterpart as would be likely in any case.



All (~ 3) events

Event	GW150914	GW151226	LVT151012
Signal-to-noise ratio	23.7	13.0	9.7
ρ			
False alarm rate FAR/yr $^{-1}$	$< 6.0 \times 10^{-7}$	$< 6.0 \times 10^{-7}$	0.37
p-value	7.5×10^{-8}	7.5×10^{-8}	0.045
Significance	$> 5.3 \sigma$	$> 5.3 \sigma$	1.7σ
Primary mass $m_1^{\text{source}}/M_\odot$	$36.2^{+5.2}_{-3.8}$	$14.2^{+8.3}_{-3.7}$	23^{+18}_{-6}
Secondary mass $m_2^{\text{source}}/M_\odot$	$29.1^{+3.7}_{-4.4}$	$7.5^{+2.3}_{-2.3}$	13^{+4}_{-5}
Chirp mass $\mathcal{M}^{\text{source}}/M_\odot$	$28.1^{+1.8}_{-1.5}$	$8.9^{+0.3}_{-0.3}$	$15.1^{+1.4}_{-1.1}$
Total mass $M^{\text{source}}/M_\odot$	$65.3^{+4.1}_{-3.4}$	$21.8^{+5.9}_{-1.7}$	37^{+13}_{-4}
Effective inspiral spin χ_{eff}	$-0.06^{+0.14}_{-0.14}$	$0.21^{+0.20}_{-0.10}$	$0.0^{+0.3}_{-0.2}$
Final mass $M_f^{\text{source}}/M_\odot$	$62.3^{+3.7}_{-3.1}$	$20.8^{+6.1}_{-1.7}$	35^{+14}_{-4}
Final spin a_f	$0.68^{+0.05}_{-0.06}$	$0.74^{+0.06}_{-0.06}$	$0.66^{+0.09}_{-0.10}$
Radiated energy $E_{\text{rad}}/(M_\odot c^2)$	$3.0^{+0.5}_{-0.4}$	$1.0^{+0.1}_{-0.2}$	$1.5^{+0.3}_{-0.4}$
Peak luminosity $\ell_{\text{peak}}/(\text{erg s}^{-1})$	$3.6^{+0.5}_{-0.4} \times 10^{56}$	$3.3^{+0.8}_{-1.6} \times 10^{56}$	$3.1^{+0.8}_{-1.8} \times 10^{56}$
Luminosity distance D_L/Mpc	420^{+150}_{-180}	440^{+180}_{-190}	1000^{+500}_{-500}
Source redshift z	$0.09^{+0.03}_{-0.04}$	$0.09^{+0.03}_{-0.04}$	$0.20^{+0.09}_{-0.09}$
Sky localization $\Delta\Omega/\text{deg}^2$	230	850	1600



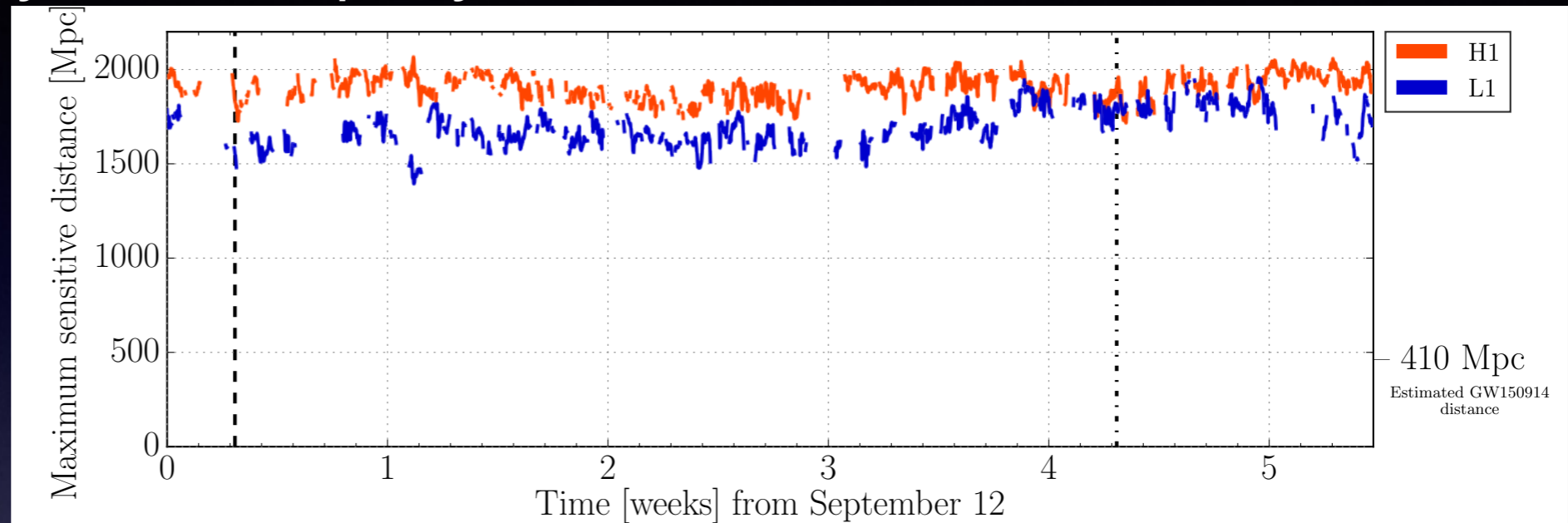
Rate of gravitational waves from BH-BH

Very simple one:

1 event in 16 live days. \rightarrow 25 per yr.

sensitivity redshift,
z of 0.3, 1.6 Gpc
 \rightarrow Vol ~ 7 Gpc³

$$3.5 \text{ Gpc}^{-3} \text{ yr}^{-1}$$



The GW was observed at high S/N, there are going to be other events (as the LVT151012). Also if BHs are from Pop III stars or are at globular clusters or at regions of low metallicity and high grav. potential they will have some mass distribution and also will have some redshift distribution.

Going over astrophysical uncertainties in the above assumptions:

Using only GW150914 (fixing the masses, spins): $2 - 53 \text{ Gpc}^{-3} \text{ yr}^{-1}$

Using both GW150914 and LVT151012: $6 - 400 \text{ Gpc}^{-3} \text{ yr}^{-1}$

LIGO's combined range: $2 - 400 \text{ Gpc}^{-3} \text{ yr}^{-1}$

LIGO's upgraded O1 (2015-16) run:

Mass distribution	$R/(\text{Gpc}^{-3}\text{yr}^{-1})$		
	PyCBC	GstLAL	Combined
Event based			
GW150914	$3.2^{+8.3}_{-2.7}$	$3.6^{+9.1}_{-3.0}$	$3.4^{+8.6}_{-2.8}$
LVT151012	$9.2^{+30.3}_{-8.5}$	$9.2^{+31.4}_{-8.5}$	$9.4^{+30.4}_{-8.7}$
GW151226	35^{+92}_{-29}	37^{+94}_{-31}	37^{+92}_{-31}
All	53^{+100}_{-40}	56^{+105}_{-42}	55^{+99}_{-41}
Astrophysical			
Flat in log mass	31^{+43}_{-21}	30^{+43}_{-21}	30^{+43}_{-21}
Power Law (-2.35)	100^{+136}_{-69}	95^{+138}_{-67}	99^{+138}_{-70}

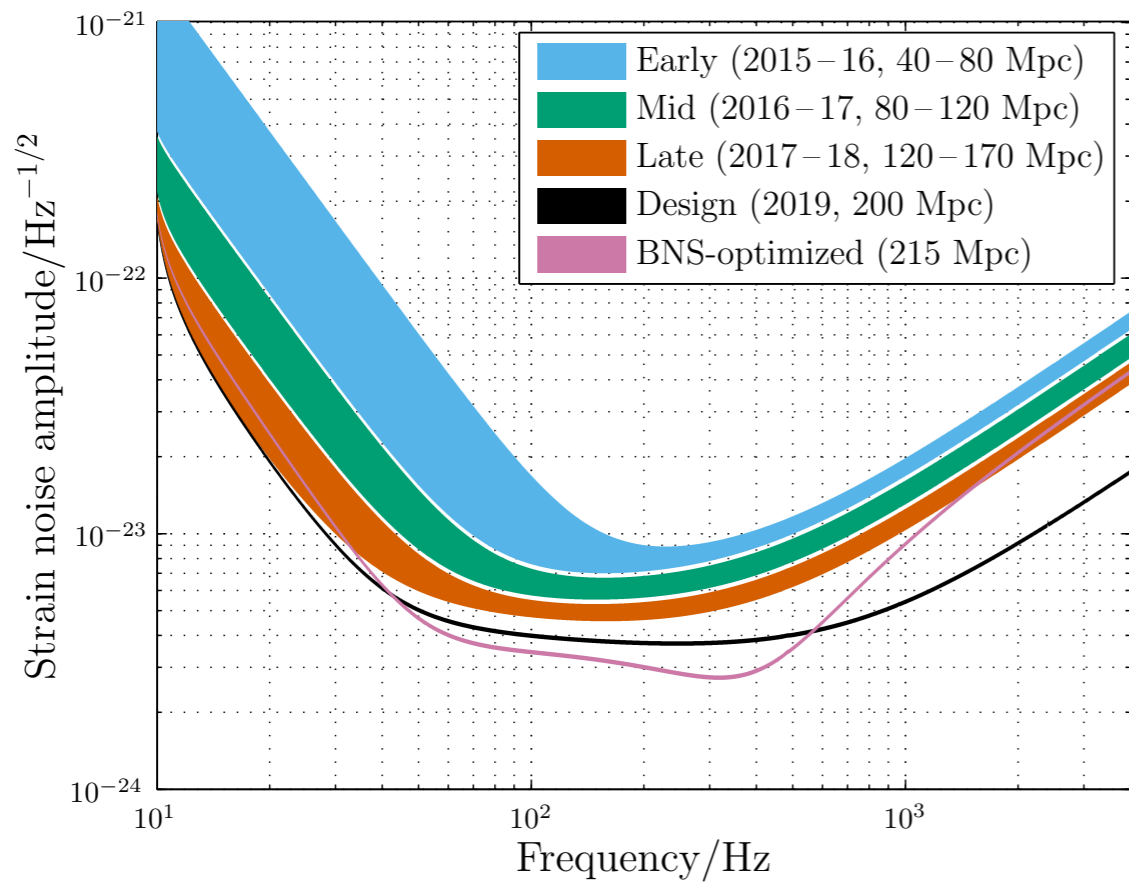
PBH?

TABLE II. Rates of BBH mergers based on populations with masses matching the observed events, and astrophysically motivated mass distributions. Rates inferred from the PyCBC and GstLAL analyses independently as well as combined rates are shown. The table shows median values with 90% credible intervals.

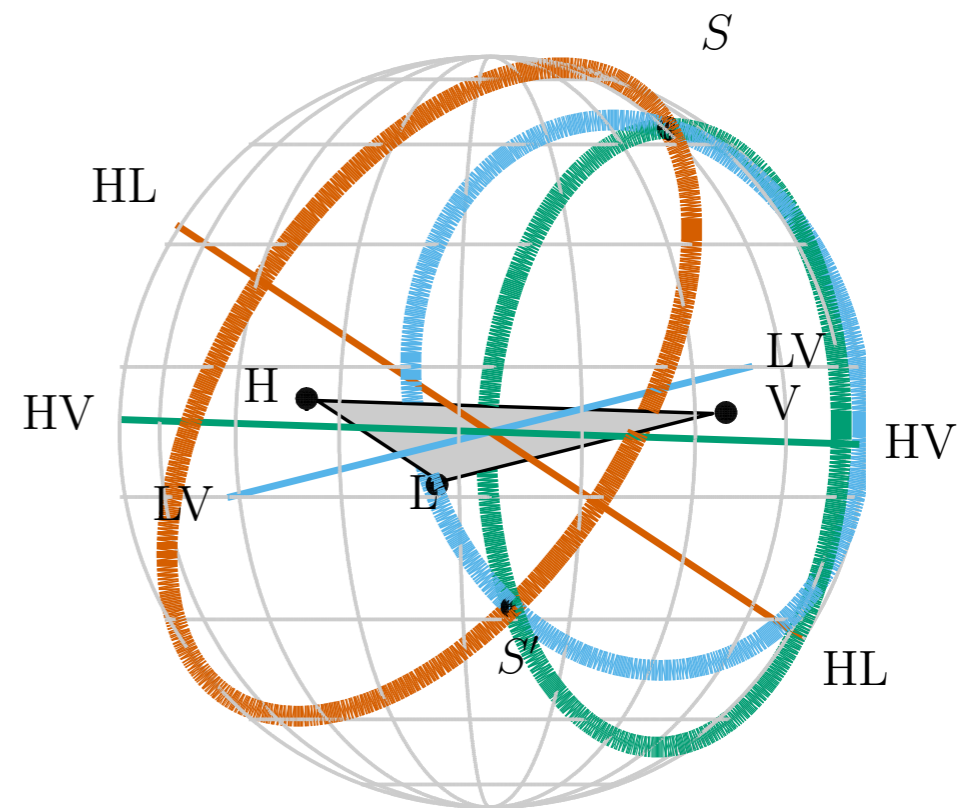
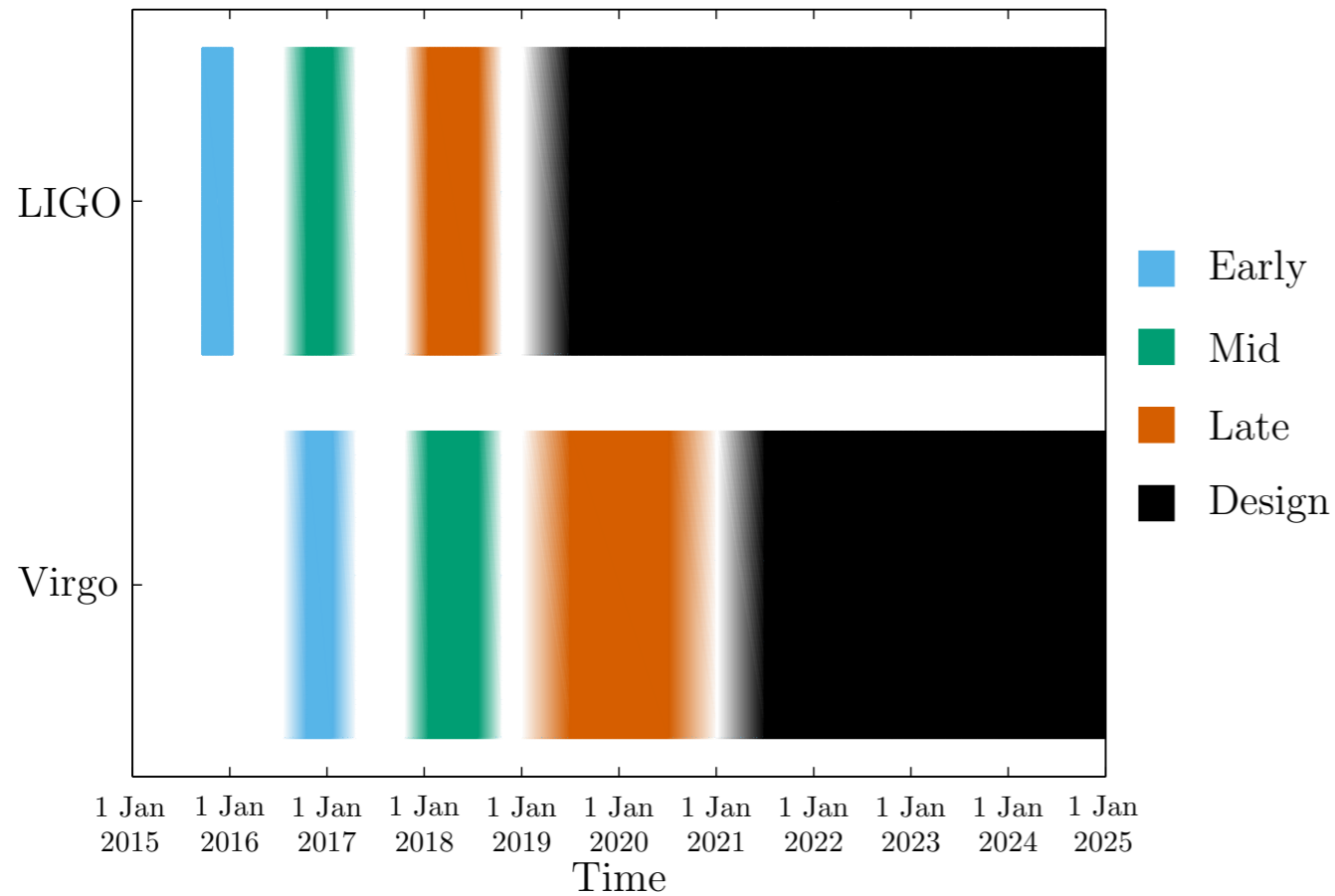
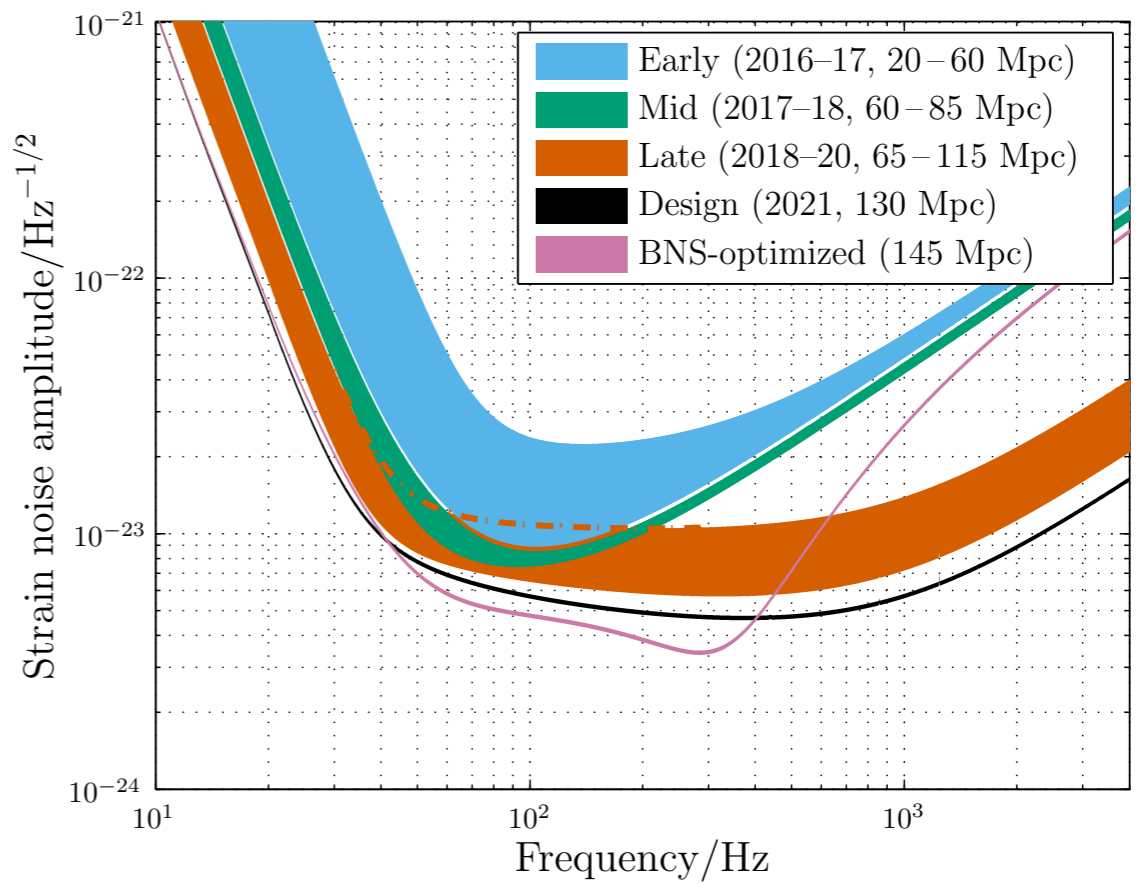
Different estimates on the coalescence rates come from different astrophysical assumptions

Future for LIGO and VIRGO

Advanced LIGO



Advanced Virgo

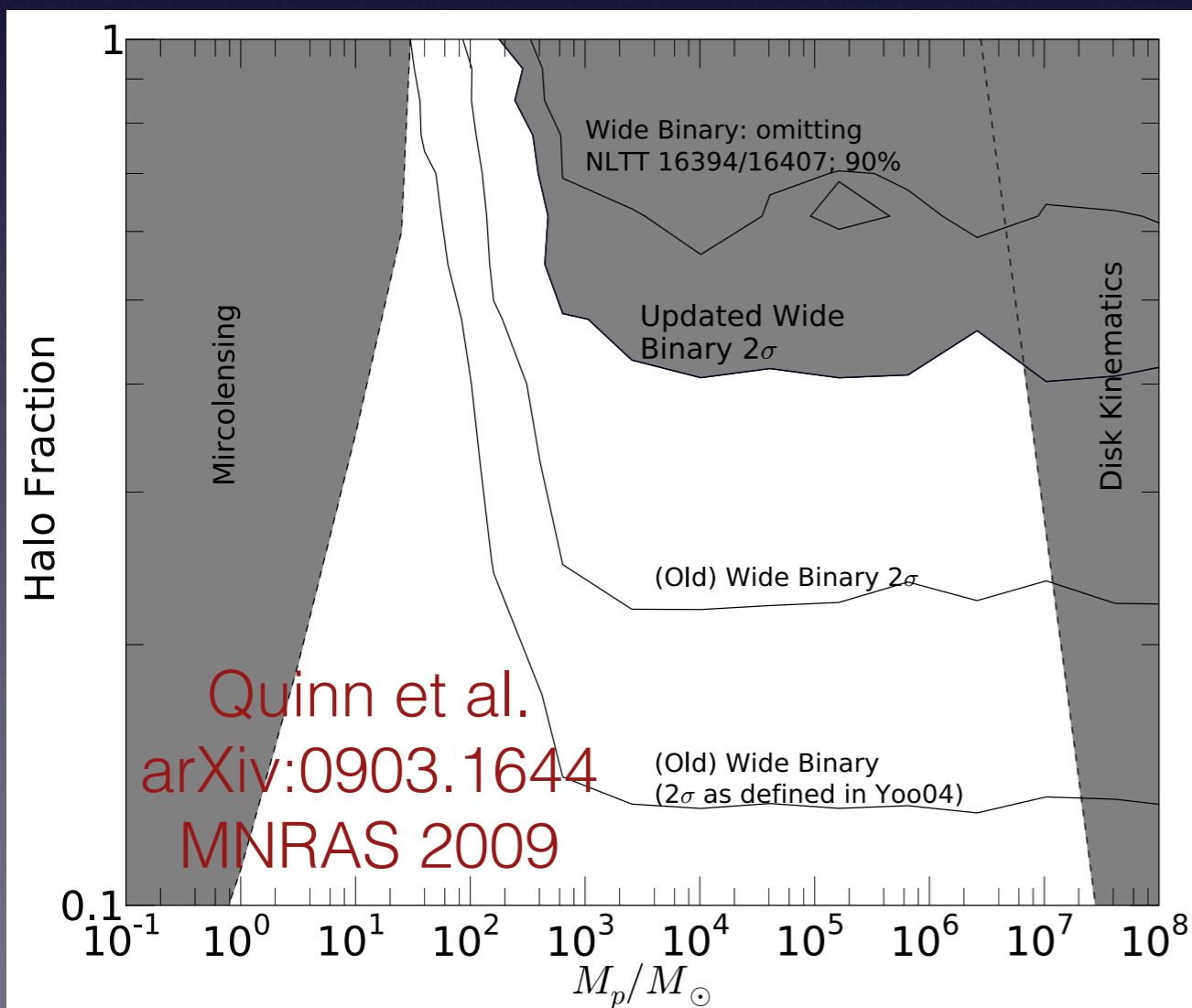


Making a connection with DM

Work with Simeon Bird, Julian B Munoz, Yacine Ali-Haimoud, Marc Kamionkowski, Ely D. Kovetz, Alvise Raccanelli and Adam Riess (JHU)
PRL 116.201031 (arXiv:1603.00464)

Assuming Dark Matter is composed by Primordial BHs.

There is some allowed parameter space around $\sim 20\text{-}70 M_{\odot}$



For the remainder I will assume that all DM is composed of PBHs and set their mass to $30 M_{\odot}$

Limits on spectral distortions of the CMB are efficient above $60 M_{\odot}$

(work by Yacine Ali-Haimoud)

Limits from GC in dwSphs (e.g. Eridanus II) (Tim Brandt arXiv:1605.03662) are robust below $15 M_{\odot}$.

Limits from micro-lensing of macro-lensed quasars depend on the DM profile and vel. dips. prof.

How fast do two BHs form a binary?

$$\sigma = 2^{3/7} \pi \left(\frac{85 \pi}{6\sqrt{2}} \right)^{2/7} R_s^2 \left(\frac{v}{c} \right)^{-18/7}$$

In easy units:

$$\sigma = 1.37 \times 10^{-14} M_{30}^2 v_{199}^{-18/7} \text{pc}^2$$

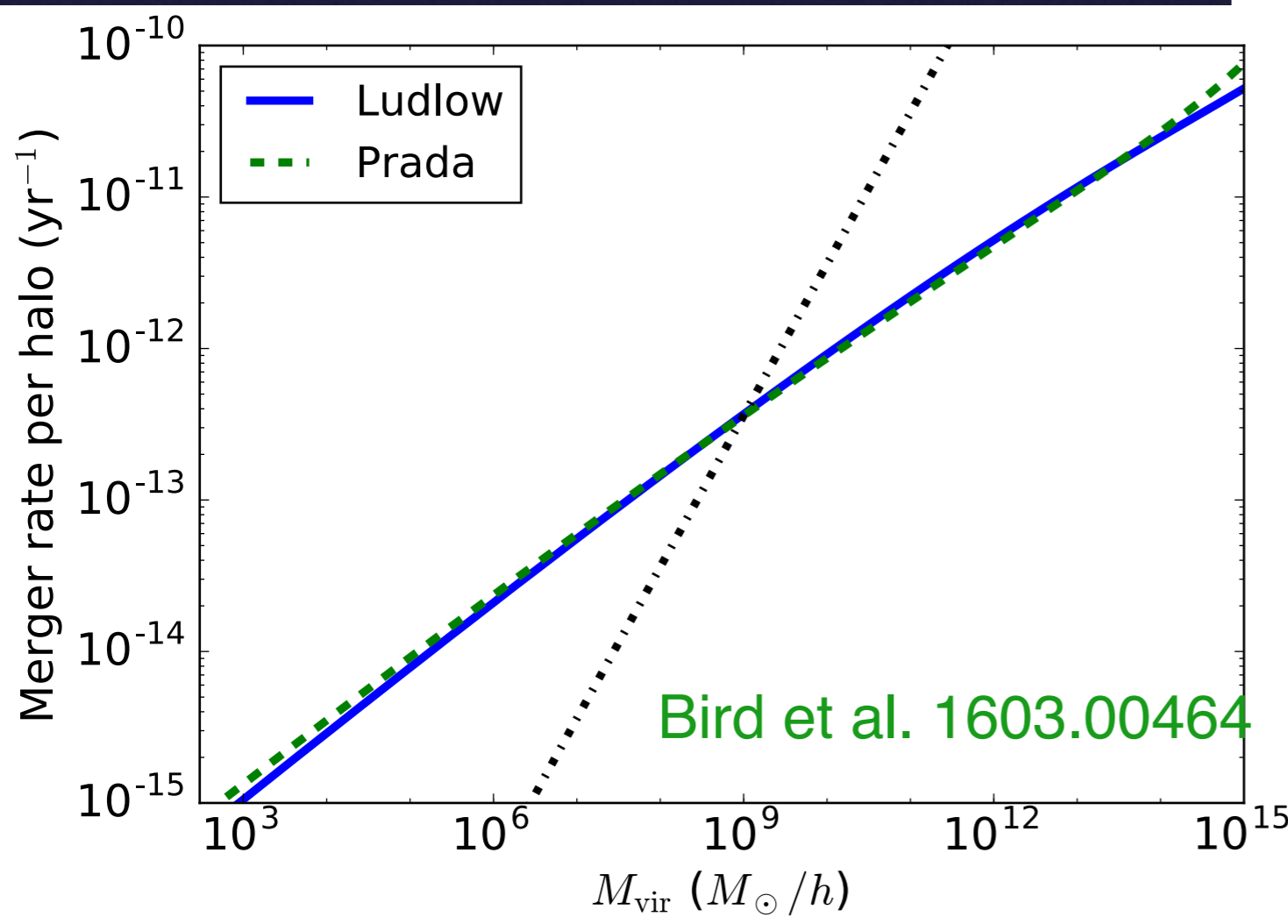
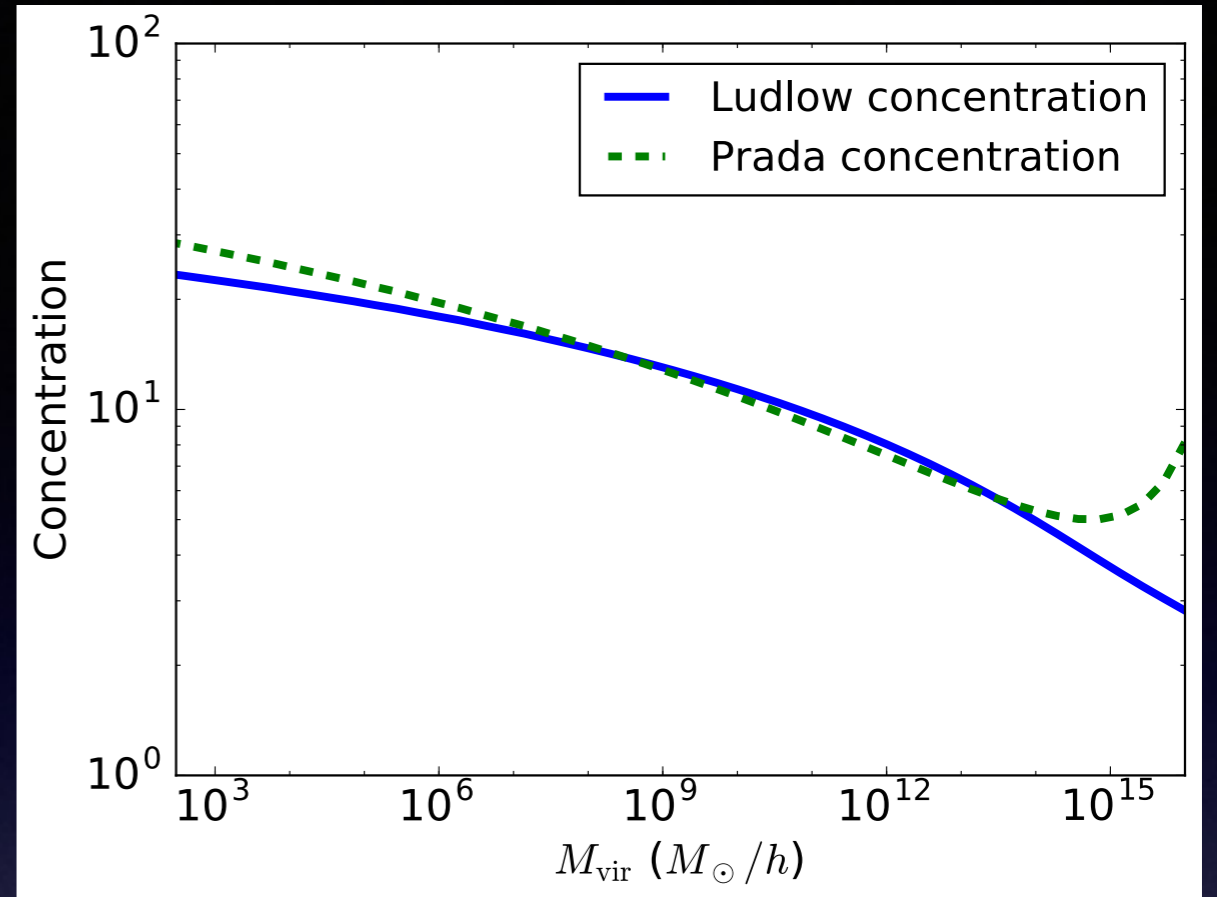
Assuming an NFW profile for the PBHs:

$$\rho_{NFW}(r) = \frac{\rho_0}{(r/R_s) \cdot (1 + r/R_s)^2}$$

One gets a Rate of PBHs mergers:

$$\mathcal{R} = 4\pi \int_0^{R_{\text{vir}}} r^2 \frac{1}{2} \left(\frac{\rho_{\text{nfw}}(r)}{M_{\text{pbh}}} \right)^2 \langle \sigma v_{\text{pbh}} \rangle dr$$

Lower mass halos \rightarrow lower velocity dispersion (i.e. higher cross-section for the binary formation) and higher concentration:

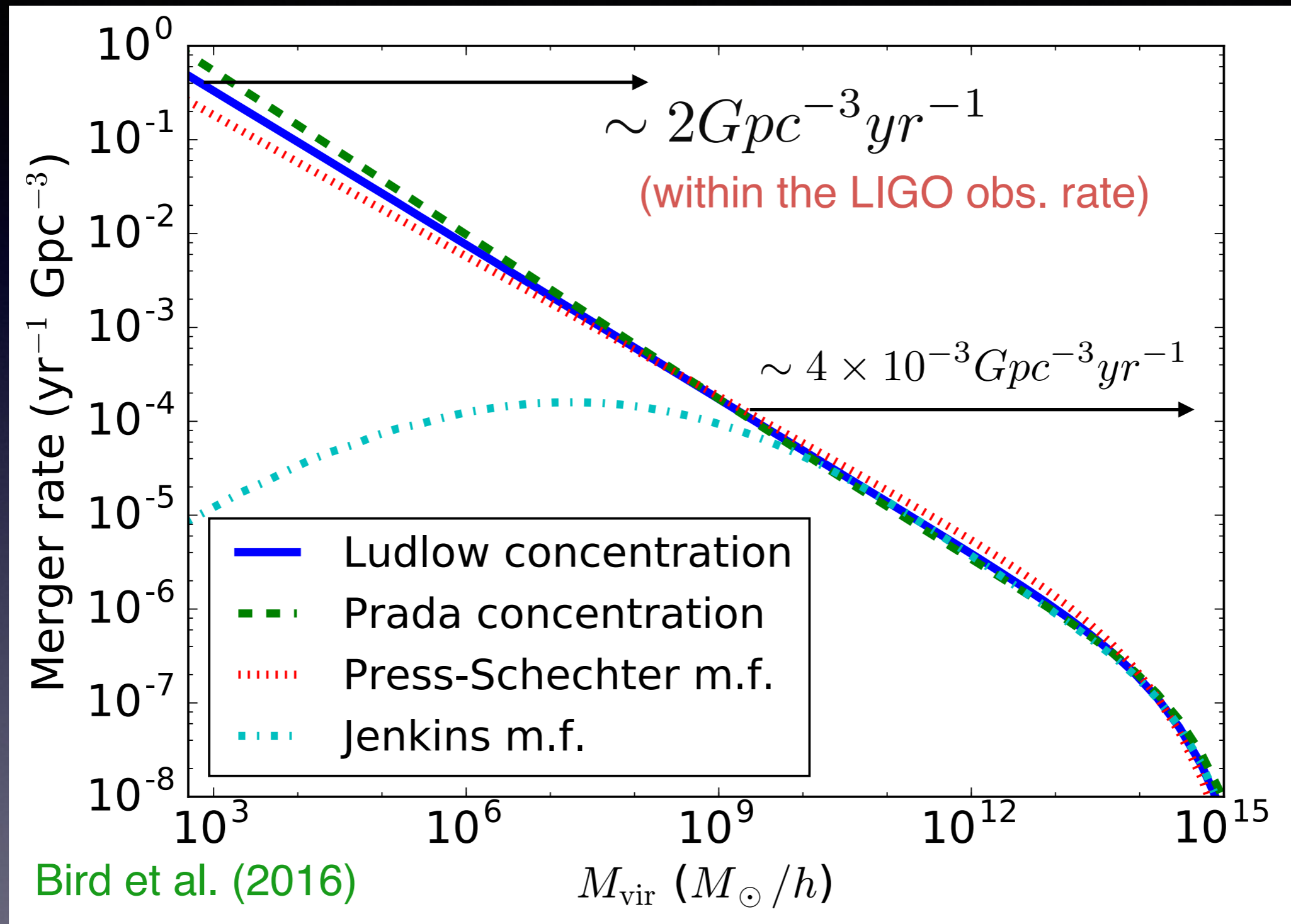


But there are many more (in terms on number) low mass DM halos:

$$\frac{dn}{dM} \sim M^{-1.85}$$

Impose a cut-off at $\sim 400 M_{\odot}$

So after including everything:



By 2019 the sensitivity will have increased to $z < 0.75$

We expect $O(10^2)$ events from PBHs (if they compose 100% of DM) by 2025.

All may be in a narrow mass range around 30 solar masses.

No other EM or neutrino signals. (typical though given that BH-BH give GW only)

Following the DM distribution (need better angular resolution though).

Basic Uncertainties in the rate calculation:

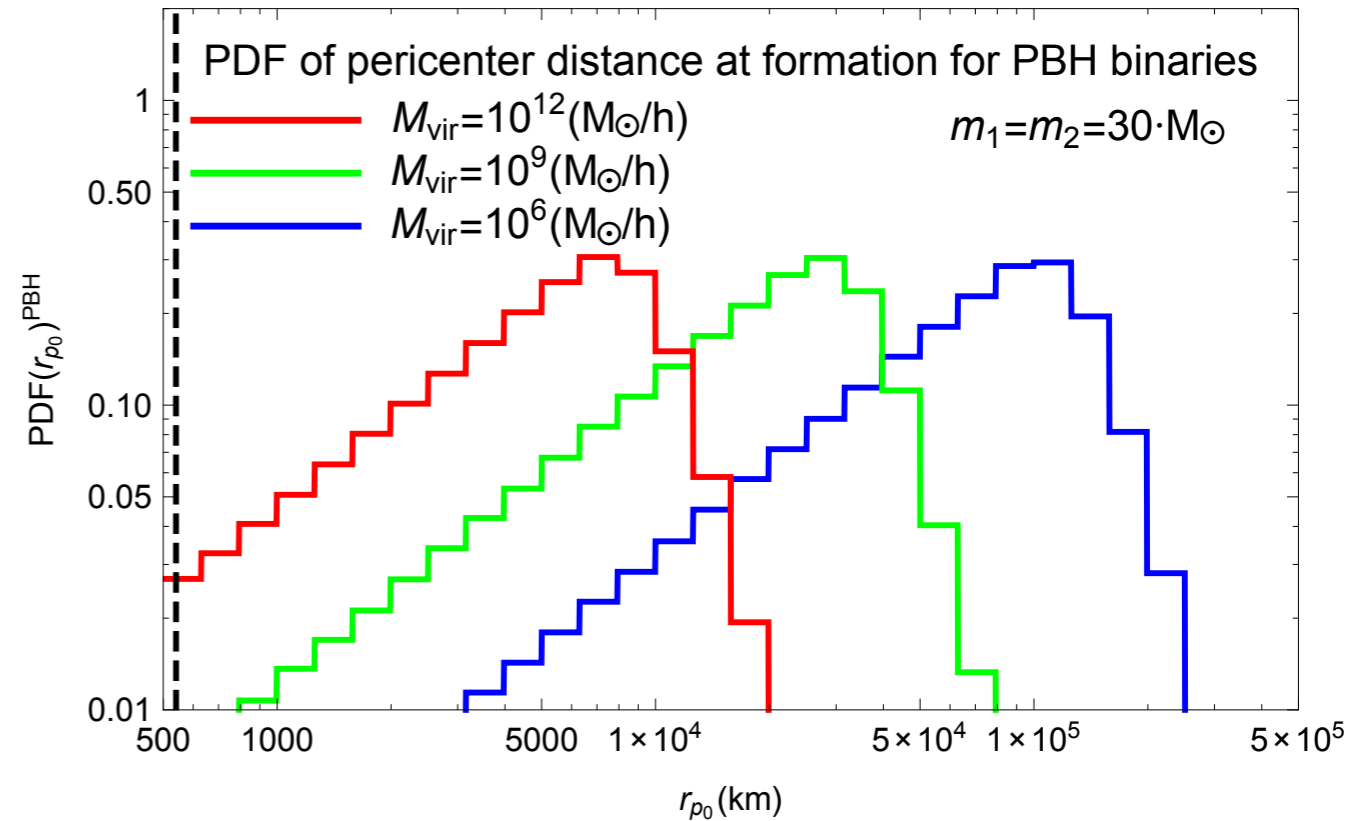
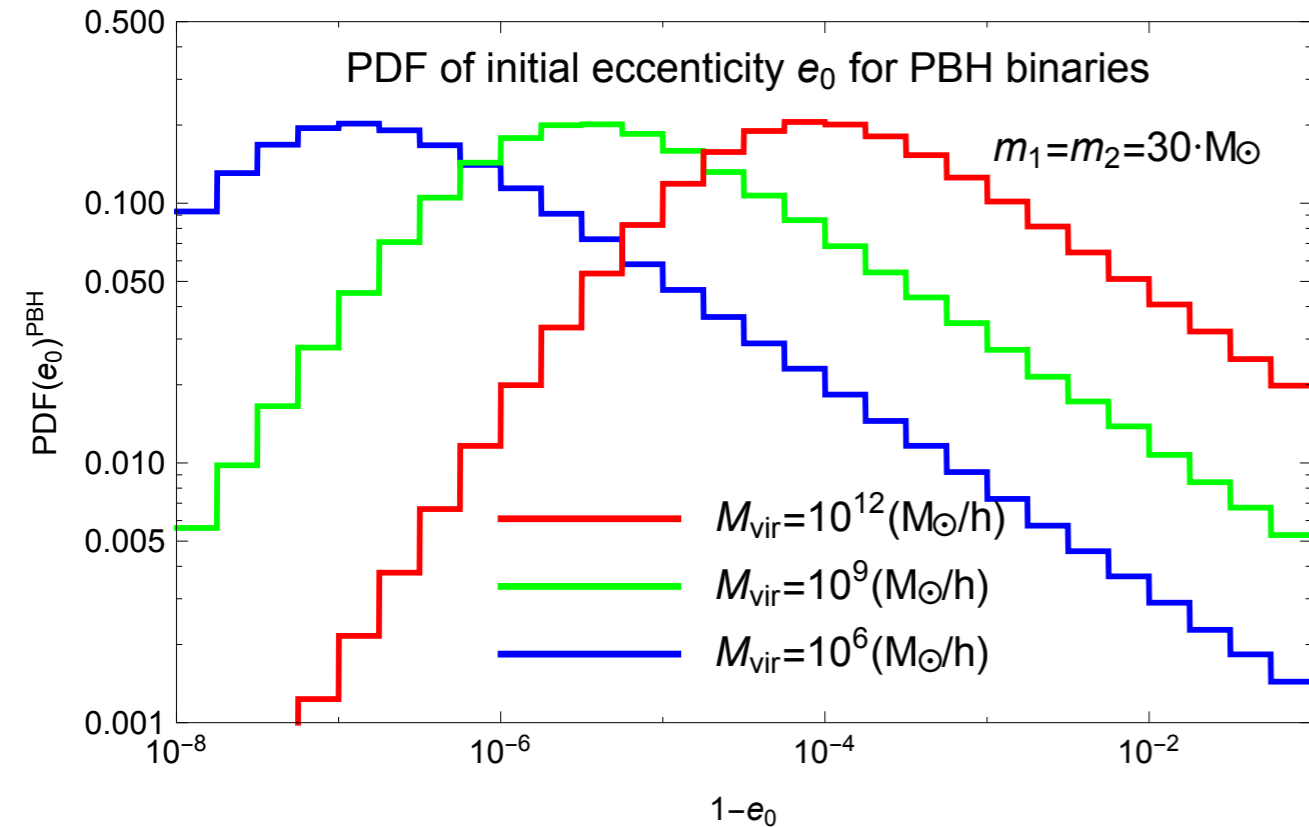
DM profile (factor of ~ 3)

Mass-Concentration relationship (factor of ~ 3)

Sub-halo contribution (previous slide) and discreteness of smallest halos.

Future directions for DM by PBHs

When these binaries form they have **high initial eccentricities** and **small peri-center distances**:



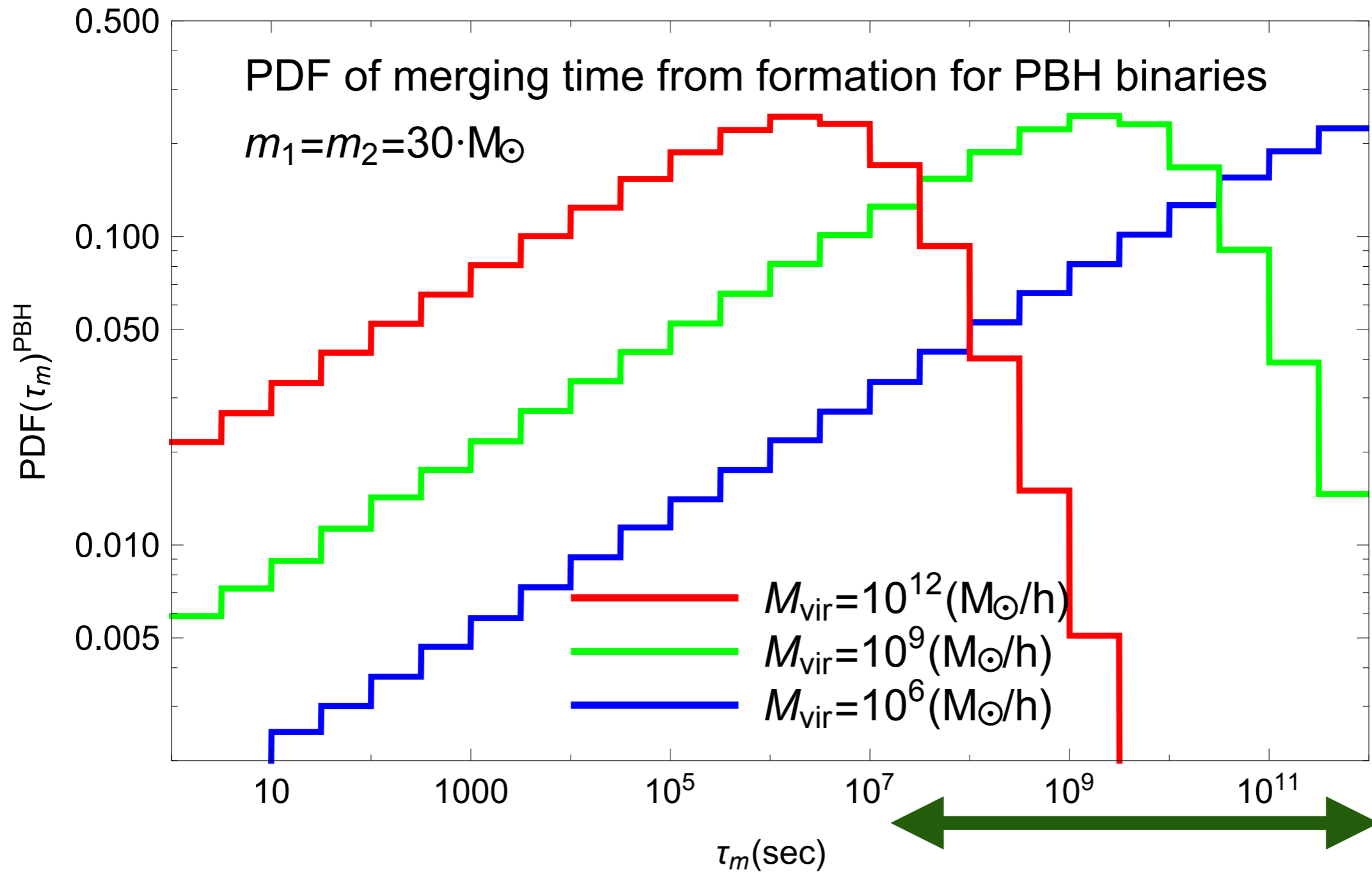
PDFs of the PBH formed binaries

$$(1 - e_0)^{\text{peak}} \simeq 2.6 \xi \eta^{2/7} (w/c)^{10/7} \quad \xi \simeq 1, \eta = 1/4 \quad \text{for equal BH masses}$$

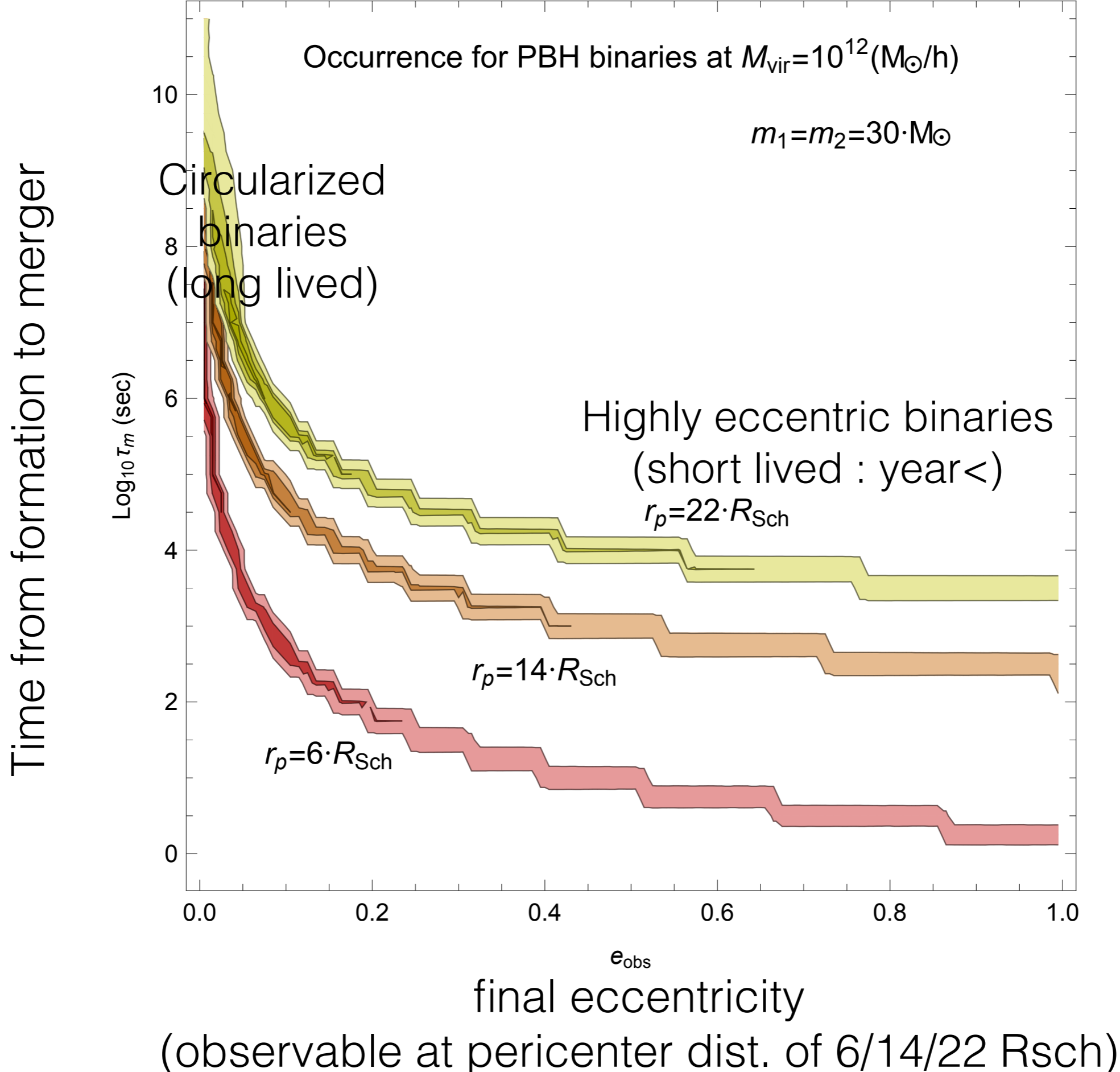
$$r_{p_0} \simeq 2 \times 10^4 \text{ km} (v_{\text{DM}}/20 \text{ km/s})^{-4/7} \quad w \simeq 2/20/200 \text{ km/s}$$

I.C., E. Kovetz, Ali-Haimoud, S. Bird, M. Kamionkowski, J. Munoz and A. Raccanelli (JHU) PRD 94 084013 (arXiv:1606.07437)

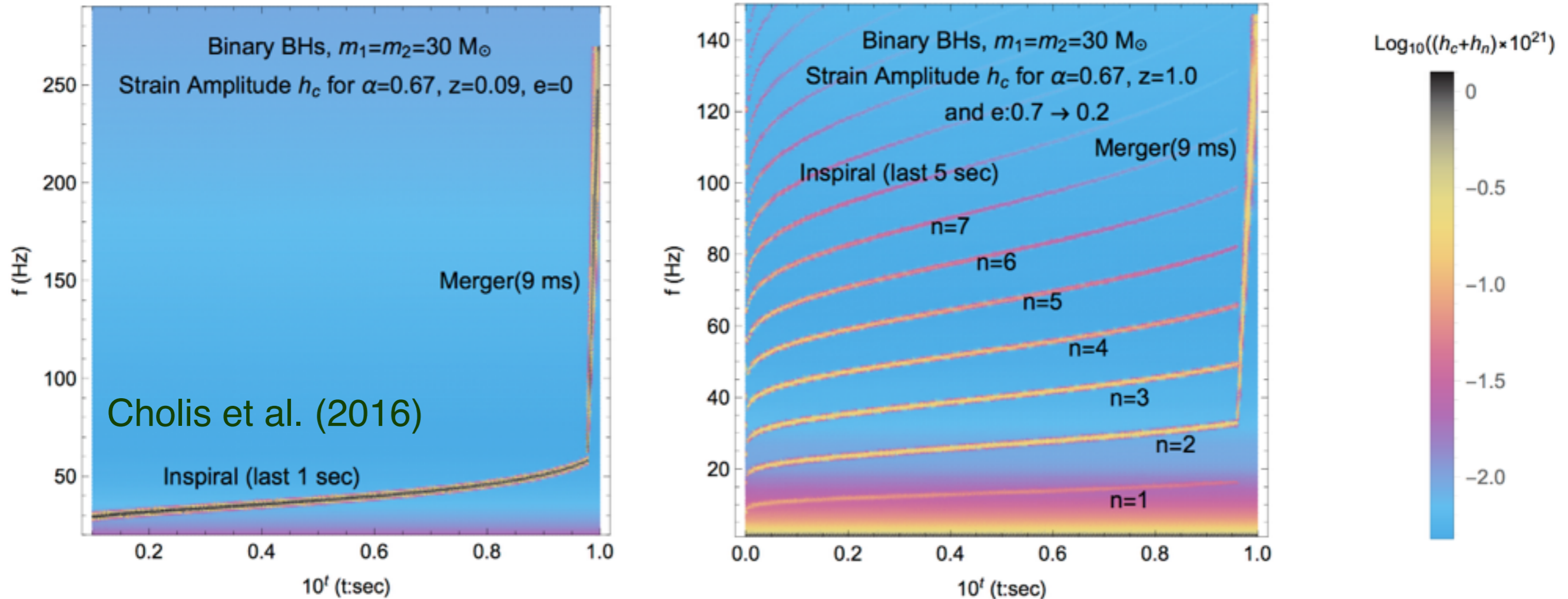
Which in turn have dramatically different timescales until merger:



By the time of LIGO observation fully circularized.



A rare case? (see many more modes of grav. waves)



simplified noise (LIGO final design)

With LIGO we expect $O(1)$ events while with the Einstein Telescope we expect $O(10)$ events with multiple modes detected from PBH binaries. Other astrophysical mechanisms for Binary BHs have typical time-scales of evolution that is \sim Myrs-Gyrs. With Future eLISA we will also be able to trace back some PBH systems to earlier stages (days-years before the merger event) and thus observe the binaries at even higher eccentricities.

Future Direction: The stochastic GW background

For every event like the GW150914 there are many more too distant or not powerful enough to be resolved above the threshold. These create a “stochastic” grav. wave background.

The energy density of GWs can be described by:

$$\Omega_{GW} = \frac{f}{\rho_c} \frac{d\rho_{GW}}{df} \leftarrow \text{energy density between } f \text{ and } f+df$$

$$\Omega_{GW} = \frac{f}{\rho_c H_0} \int_0^{z_{max}} dz \frac{R_m(z, \theta_k)}{(1+z) \sqrt{\Omega_\Lambda + \Omega_M (1+z)^3}} \frac{dE_{GW}(f_s, \theta_k)}{df_s}$$

f_s :frequency at source

θ_k :astrophysics assumptions,
(mass distr. of BHs and z-distr.)

energy density
spectrum, for
inspiral typically
 $\sim M_c^{5/3} f_s^{-1/3}$

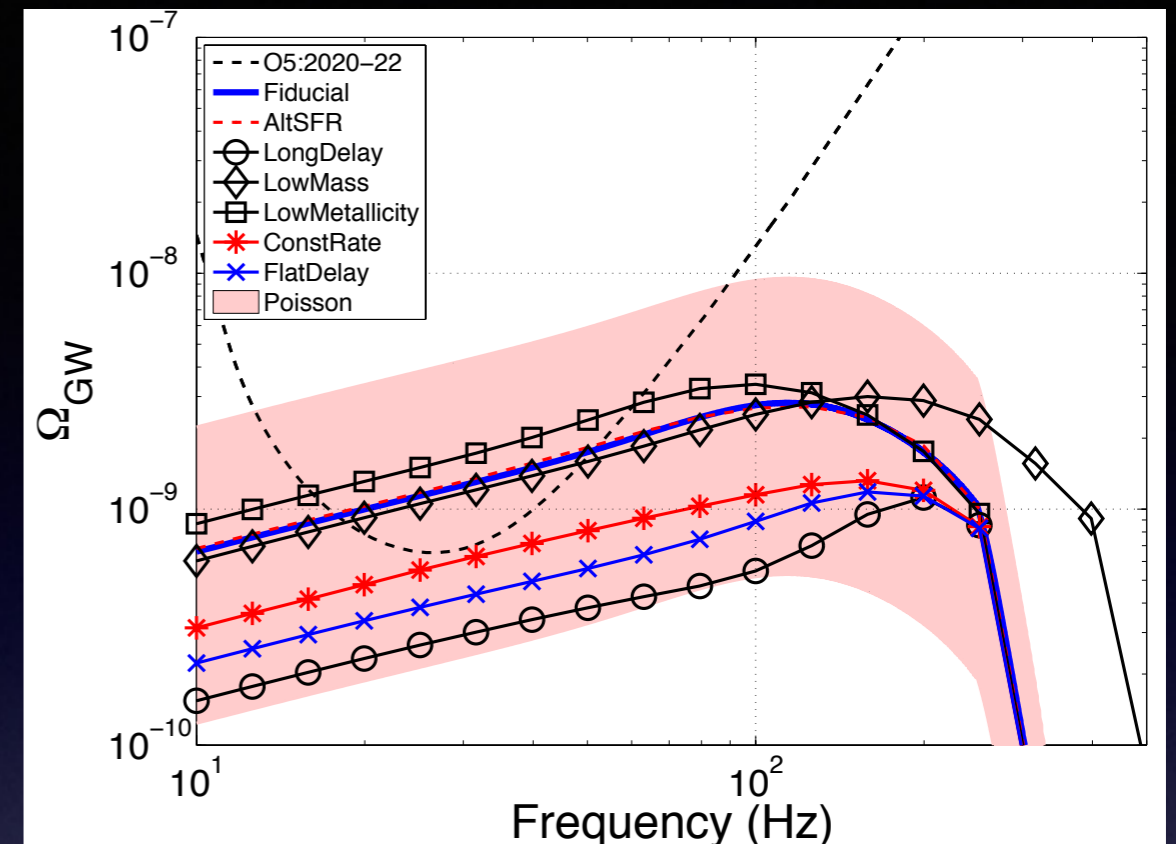
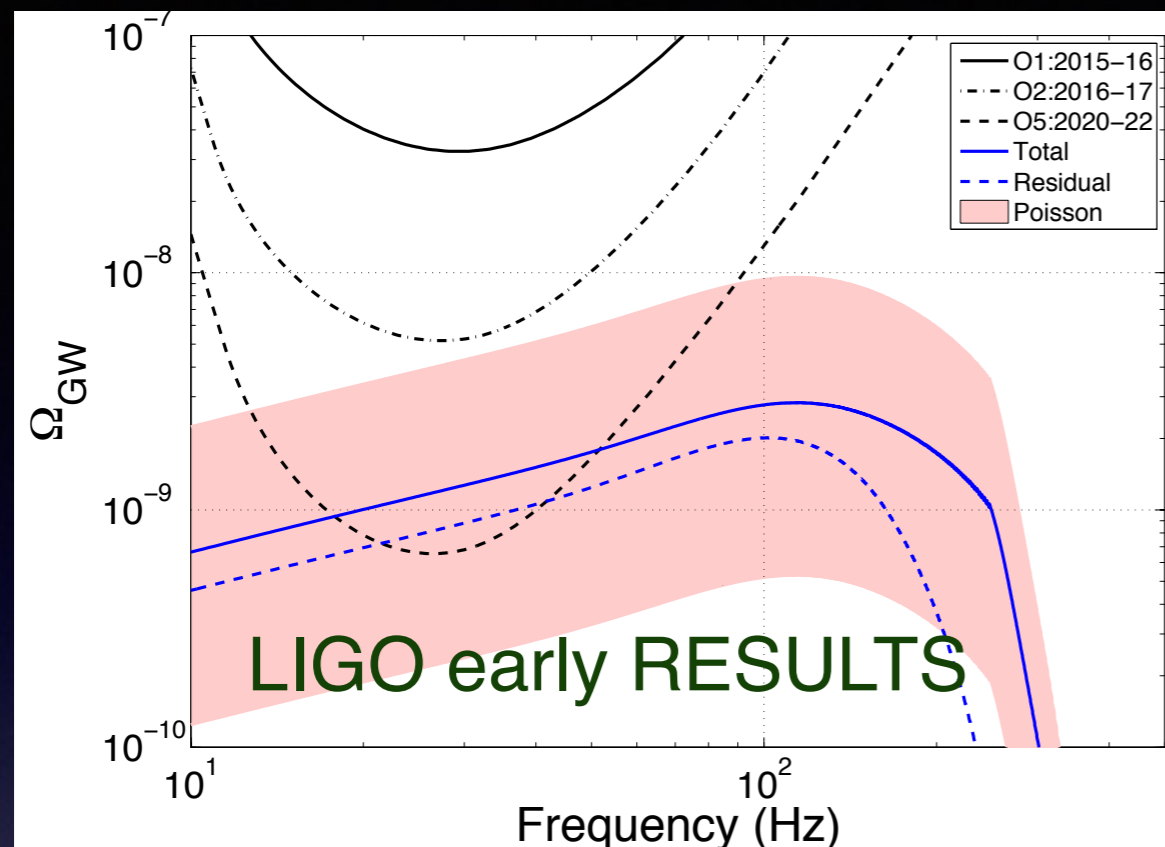
$$R_m(z, \theta_k) = \int_{t_{min}}^{t_{max}} R_f(z_f, \theta_k) P(t_d, \theta_k) dt_d$$

distr. of time delay

:rate of BH-BH merger

binary formation rate

Measuring the stock. back will probe the GW sources



Based on the rate of $2 - 53 \text{ Gpc}^{-3} \text{ yr}^{-1}$ and assuming a conventional

Star Formation Rate (SFR) “Fiducial”

Star Formation Rate doesn’t affect much such a calculation (“AltSFR”)

“Long Delay”: it takes at least 5 Gyrs for a merger to occur (largely separated objects with slow rel. velocity before binary creation). “Flat delay” : 1 Gyr.

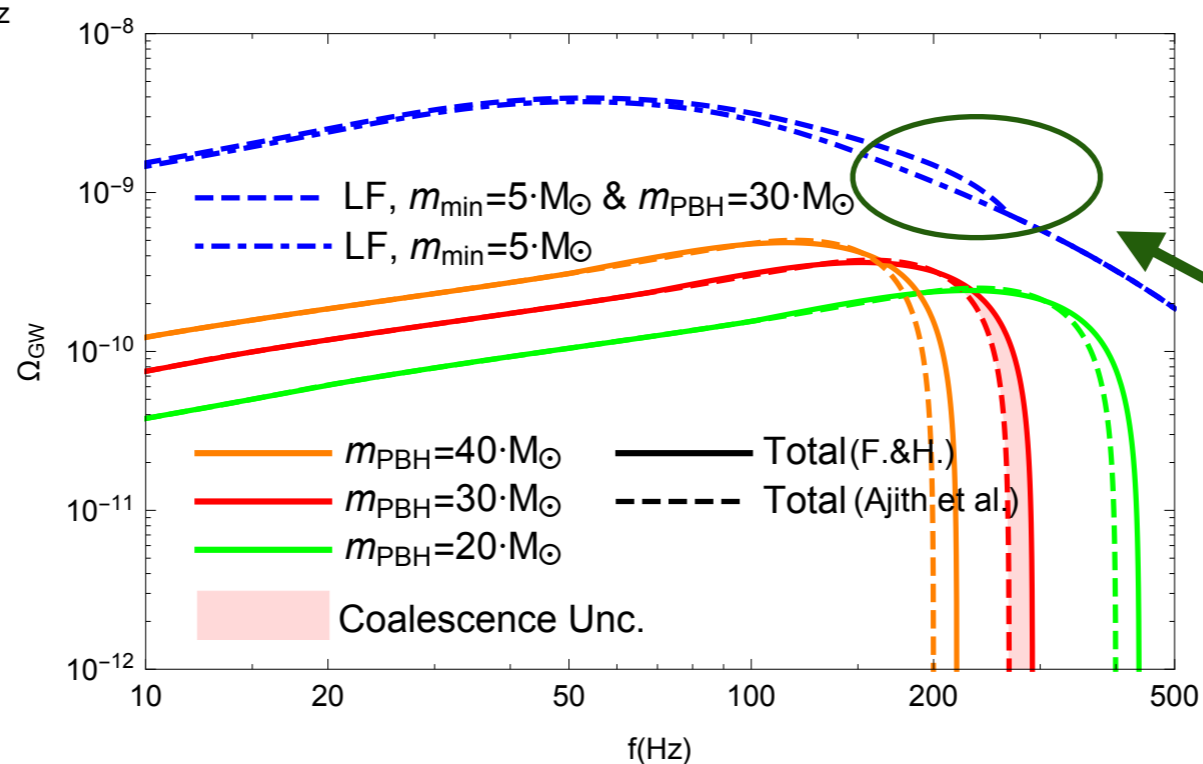
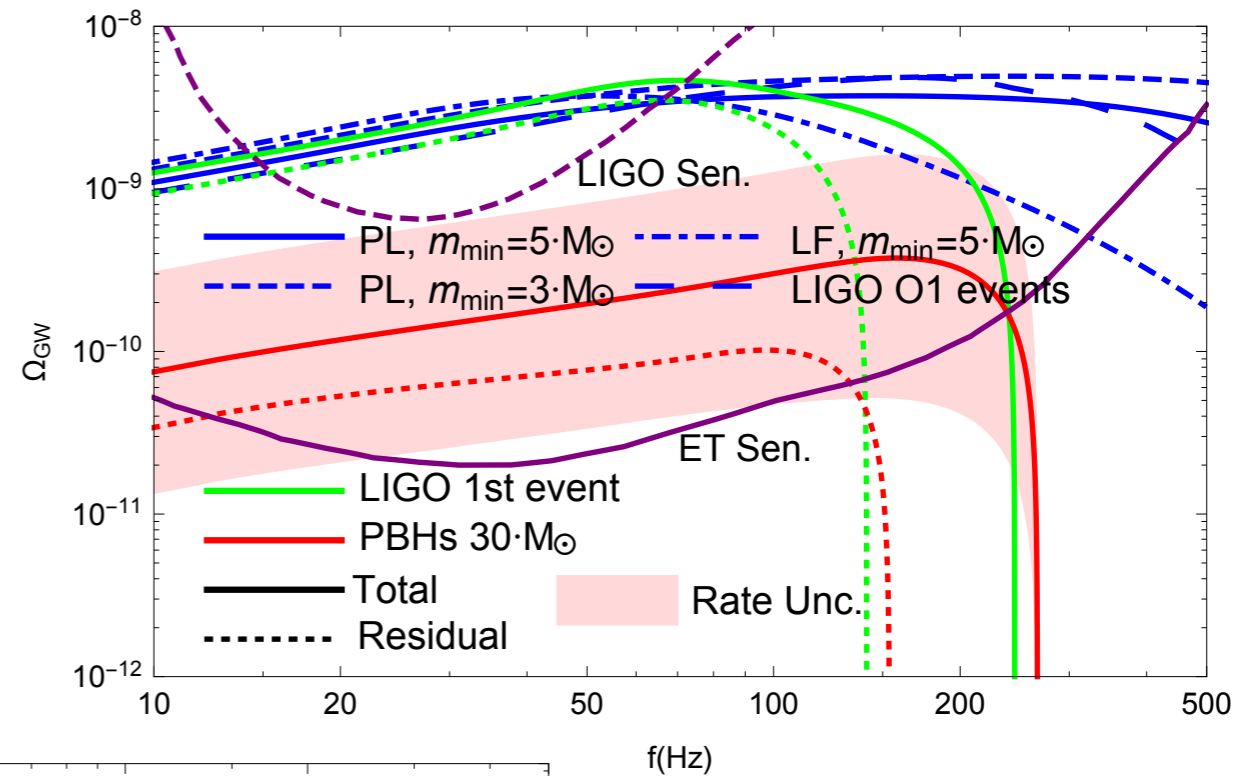
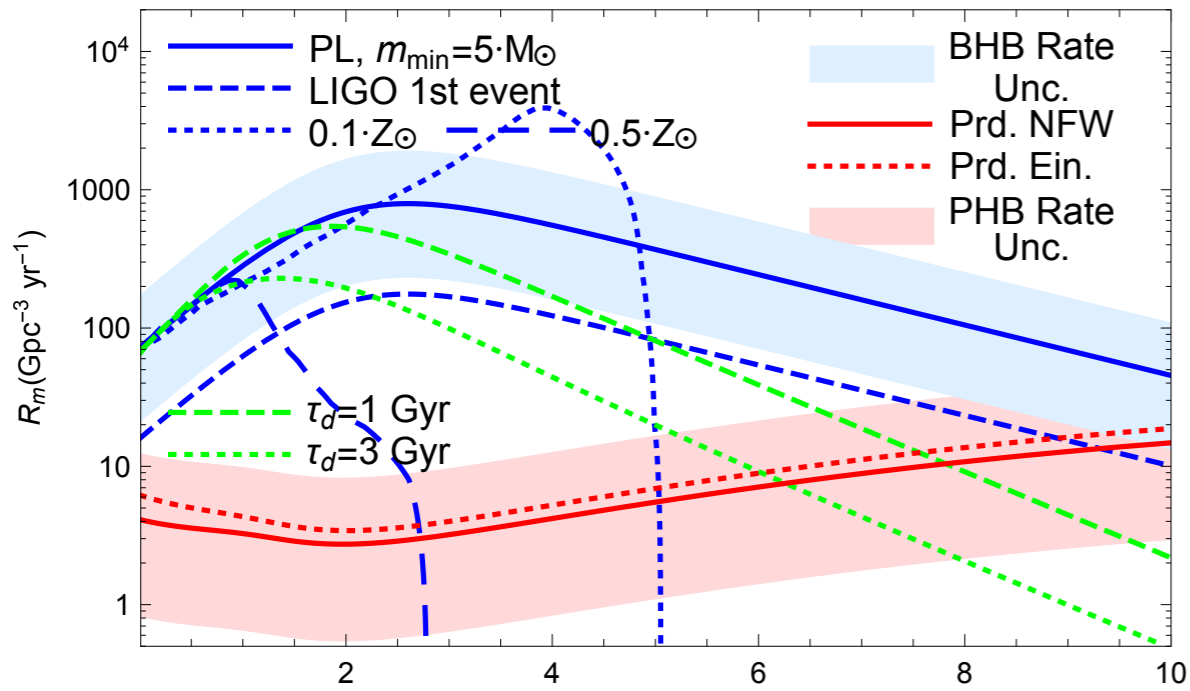
“Low mass”: assuming $15 M_{\odot}$ BHs. More power at higher frequencies.

Lower metallicity increases the number density of BHs

“Constant (in z) rate”: $R_m(z) = 16 \text{ Gpc}^{-3} \text{ yr}^{-1}$

Updated Rates on the BH-BH mergers (some room a PBH component to be seen in the Stoch. Background)

V. Mandic, S. Bird, I.C. (PRL accept.) arXiv:1608.06699 & I.C. arXiv:1609.03565



With Einstein Telescope we might be able to probe the PBH model

An other future direction: Cross-Correlations with Galaxies

A. Raccanelli, E. Kovetz, S. Bird, I.C. J. Munoz
PRD 94 023516 (arXiv:1605:01405)

If the GW signal comes from BHs originating by standard astrophysical sources e.g. BH in globular clusters, then **the binary systems should preferentially reside in galaxies where most of the stars are**. So GW and star forming galaxy (SFG) maps would be highly correlated.

If the BH binaries are mostly populating halos with different mass range, bias, redshift and angular distributions, then the correlation with SFGs galaxies in halos of masses $\sim 10^{11} - 10^{12} M_{\odot}$ would be lower.

If the GW signal comes from PBHs that constitute the DM then their distribution will be **more uniform** on the sky.

We can calculate angular projections:

$$C_{\ell}^{XY} = \langle a_{\ell m}^X a_{\ell m}^{Y*} \rangle = 4\pi \int \frac{dk}{k} \Delta^2(k) W_{\ell}^X(k) W_{\ell}^Y(k)$$

Window functions

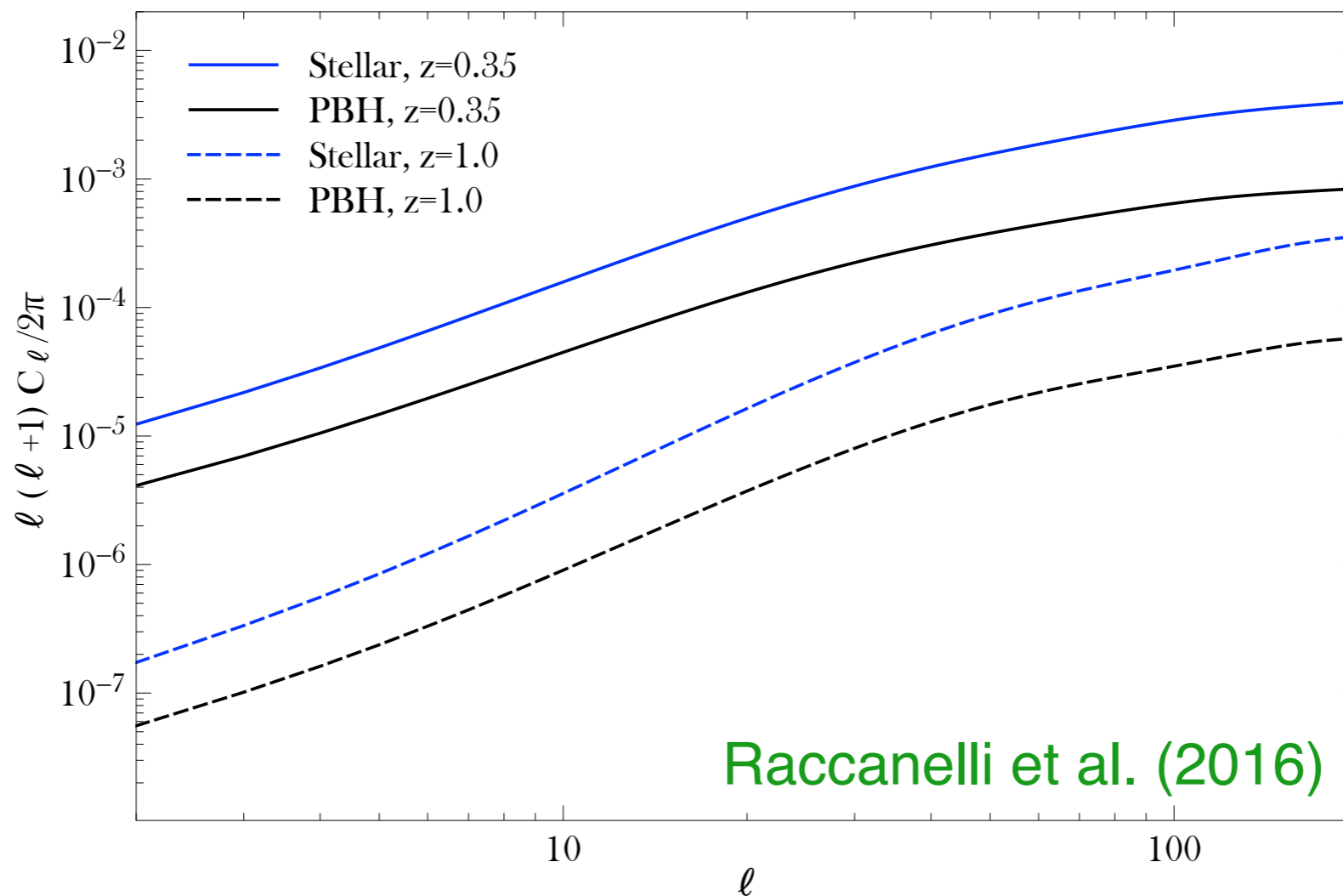
Window function:

$$W_{\ell}^X(k) = \int \frac{N_X(z)}{\text{\#/sr}} b_X^{\leftarrow}(z) j_{\ell}[k\chi(z)] dz$$

bias (progenitor infor.)

co-moving distance

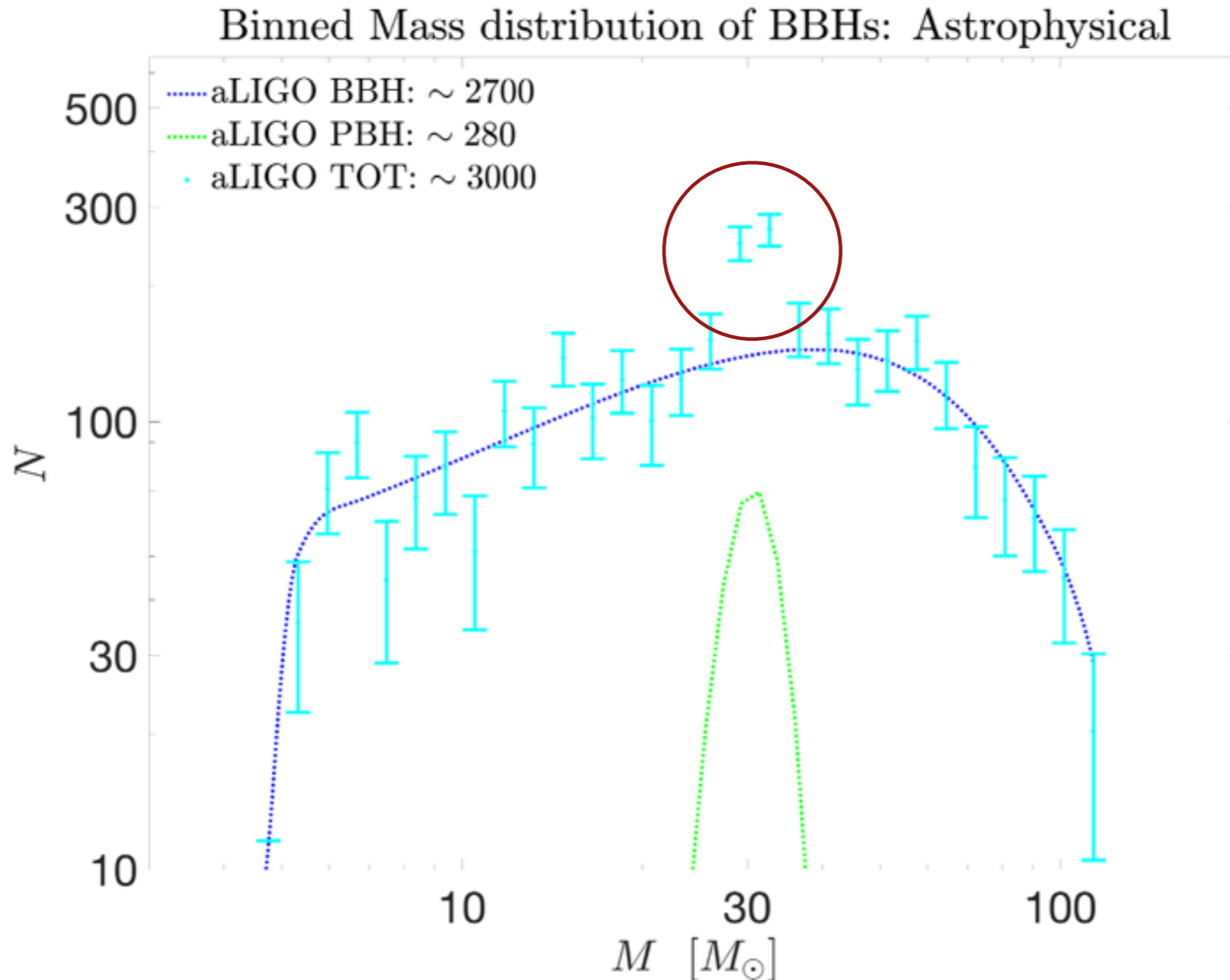
$$N_{GW}(z) = \dot{n}_{GW}(z) T_{\text{obs}} V(z)$$



Forecasted Cross-correlation amplitude of of Galaxies with BH-BH mergers. PBH binaries have a smaller bias b (~ 0.5) compared to stellar BHs (since the PBH rate is dominated by the smallest DM halos)

An other future direction: Mass-Spectrum of BH-BH binaries

E. Kovetz, I.C., P. Breysse, M. Kamionkowski arXiv:1611:01157



Conclusions: LIGO 1st Detection

Sept. 14th LIGO observed the gravitational wave signal from the coalescence of two Black Holes (GW150914)

$$m_1 = 36_{-4}^{+5} M_{\odot}$$

$$m_2 = 29_{-4}^{+4} M_{\odot}$$

$$m_{final} = 62_{-4}^{+4} M_{\odot}$$

$$m_{loss} = 3.0_{-0.5}^{+0.5} M_{\odot}$$

$$\alpha = 0.67_{-0.04}^{+0.05}$$

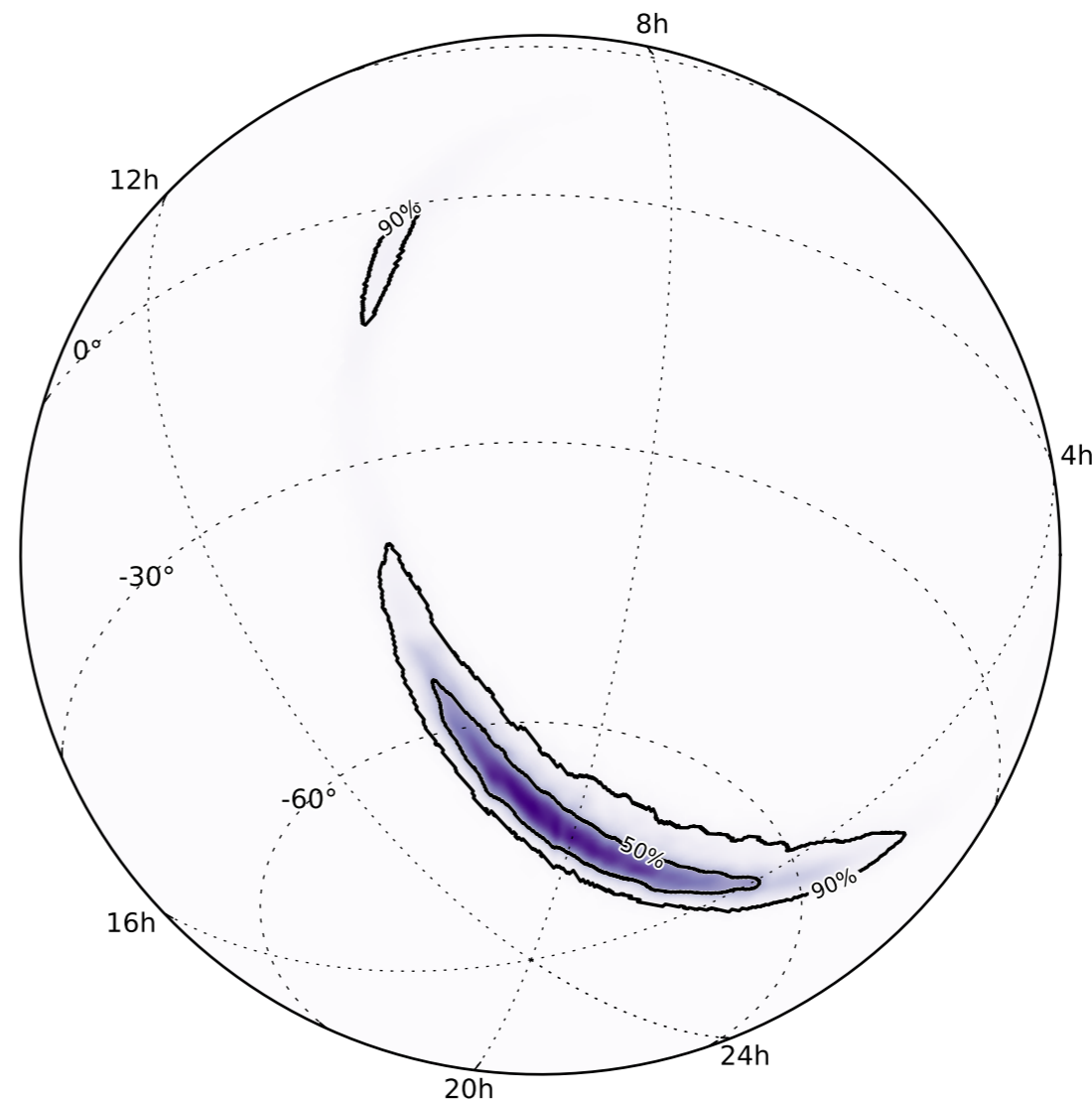
$$d_L = 410_{-180}^{+160} Mpc$$

$$z_s = 0.09_{-0.04}^{+0.03}$$

$$t_d = 6.9_{-0.4}^{+0.5} ms$$

Combined Signal to Noise : 24

Detection Significance: 5.1σ



90% probability within 590 deg^2

Rate (GW150914): $0.5 - 12 Gpc^{-3} yr^{-1}$

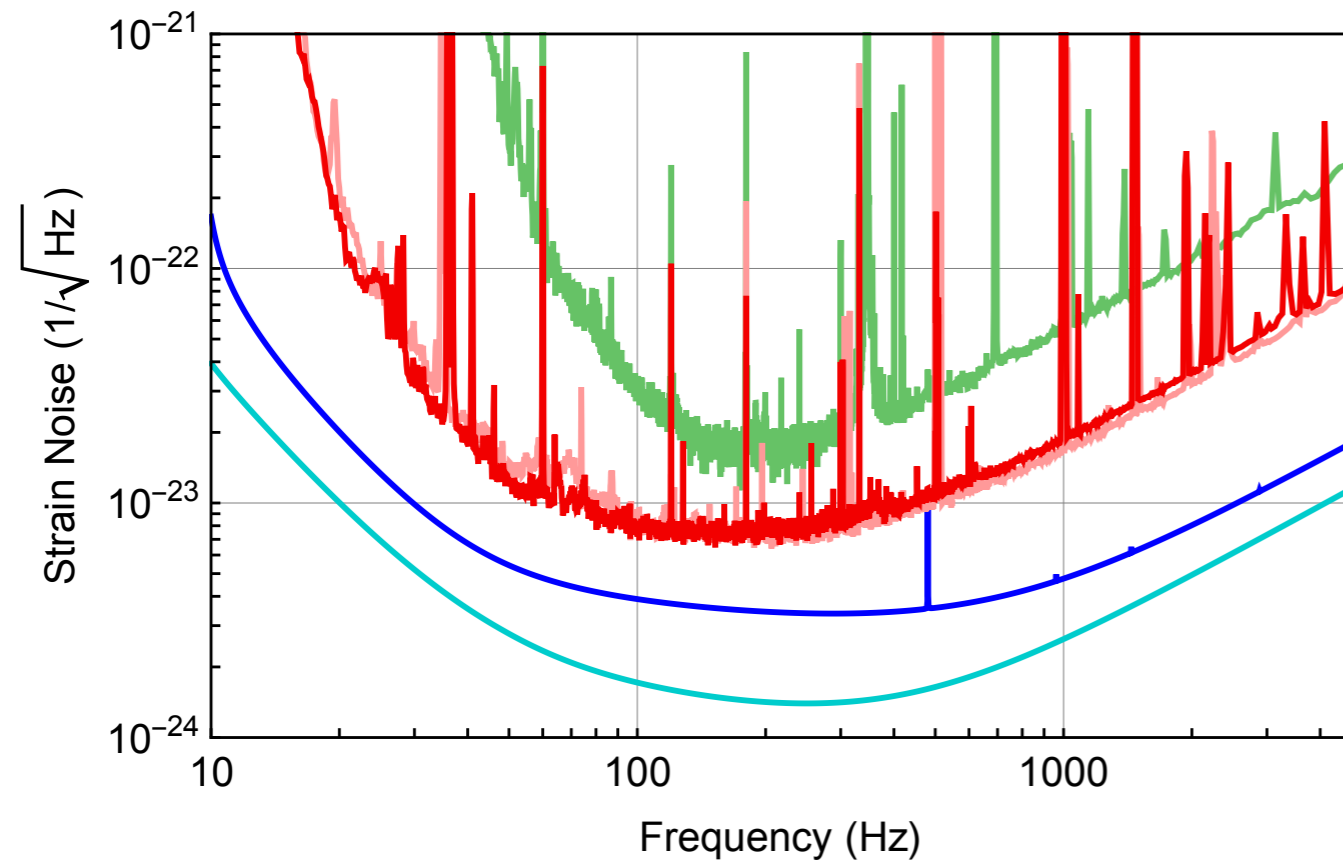
All three events: $10 - 150 Gpc^{-3} yr^{-1}$

Additional slides

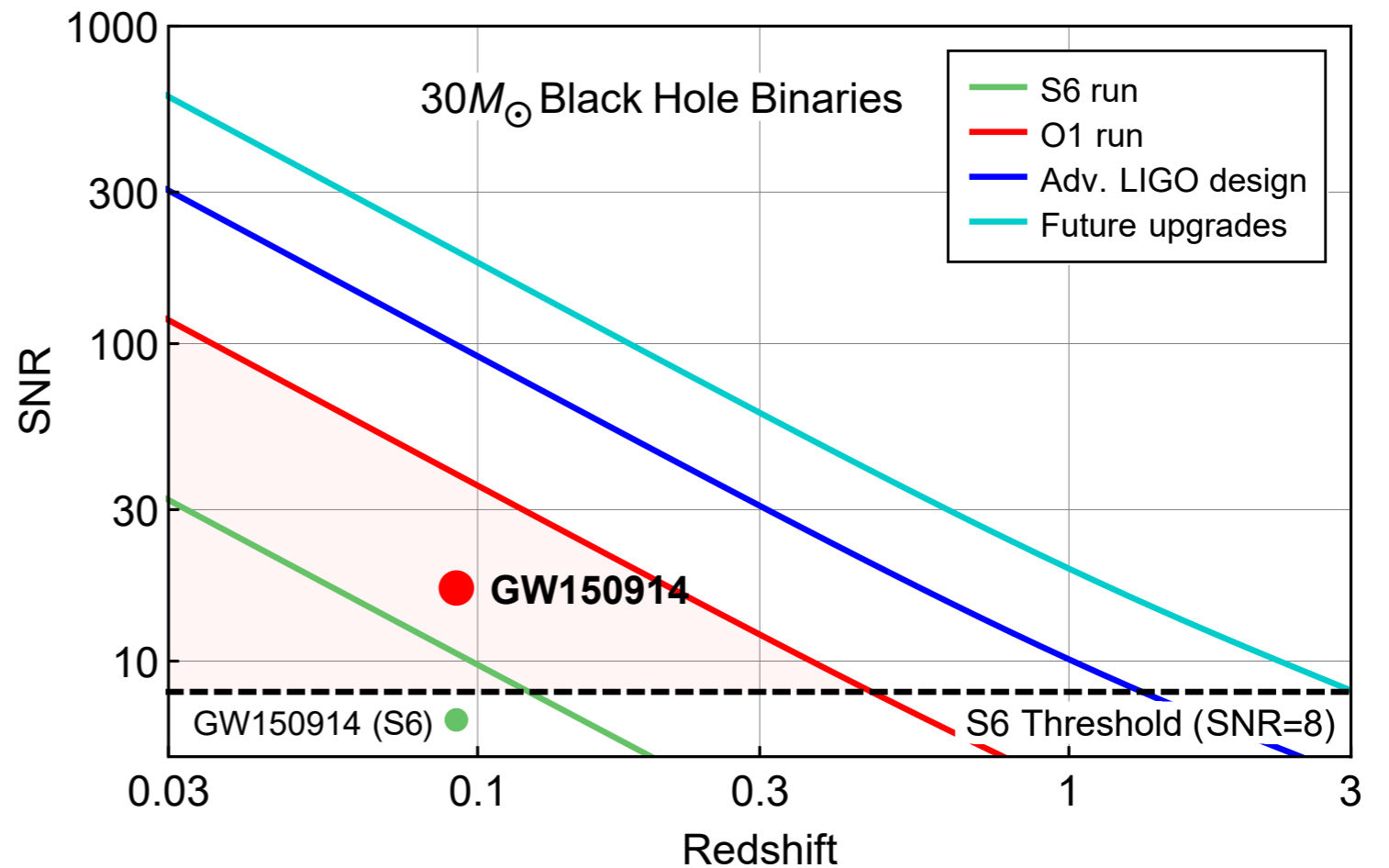
extracted parameters based on two waveform models (relying on Num. Rel. simulations)

	EOBNR	IMRPhenom	Overall
Detector-frame total mass M/M_\odot	$70.3^{+5.3}_{-4.8}$	$70.7^{+3.8}_{-4.0}$	$70.5^{+4.6\pm 0.9}_{-4.5\pm 1.0}$
Detector-frame chirp mass \mathcal{M}/M_\odot	$30.2^{+2.5}_{-1.9}$	$30.5^{+1.7}_{-1.8}$	$30.3^{+2.1\pm 0.4}_{-1.9\pm 0.4}$
Detector-frame primary mass m_1/M_\odot	$39.4^{+5.5}_{-4.9}$	$38.3^{+5.5}_{-3.5}$	$38.8^{+5.6\pm 0.9}_{-4.1\pm 0.3}$
Detector-frame secondary mass m_2/M_\odot	$30.9^{+4.8}_{-4.4}$	$32.2^{+3.6}_{-5.0}$	$31.6^{+4.2\pm 0.1}_{-4.9\pm 0.6}$
Detector-frame final mass M_f/M_\odot	$67.1^{+4.6}_{-4.4}$	$67.4^{+3.4}_{-3.6}$	$67.3^{+4.1\pm 0.8}_{-4.0\pm 0.9}$
Source-frame total mass $M^{\text{source}}/M_\odot$	$65.0^{+5.0}_{-4.4}$	$64.6^{+4.1}_{-3.5}$	$64.8^{+4.6\pm 1.0}_{-3.9\pm 0.5}$
Source-frame chirp mass $\mathcal{M}^{\text{source}}/M_\odot$	$27.9^{+2.3}_{-1.8}$	$27.9^{+1.8}_{-1.6}$	$27.9^{+2.1\pm 0.4}_{-1.7\pm 0.2}$
Source-frame primary mass $m_1^{\text{source}}/M_\odot$	$36.3^{+5.3}_{-4.5}$	$35.1^{+5.2}_{-3.3}$	$35.7^{+5.4\pm 1.1}_{-3.8\pm 0.0}$
Source-frame secondary mass $m_2^{\text{source}}/M_\odot$	$28.6^{+4.4}_{-4.2}$	$29.5^{+3.3}_{-4.5}$	$29.1^{+3.8\pm 0.2}_{-4.4\pm 0.5}$
Source-frame final mass $M_f^{\text{source}}/M_\odot$	$62.0^{+4.4}_{-4.0}$	$61.6^{+3.7}_{-3.1}$	$61.8^{+4.2\pm 0.9}_{-3.5\pm 0.4}$
Mass ratio q	$0.79^{+0.18}_{-0.19}$	$0.84^{+0.14}_{-0.21}$	$0.82^{+0.16\pm 0.01}_{-0.21\pm 0.03}$
Effective inspiral spin parameter χ_{eff}	$-0.09^{+0.19}_{-0.17}$	$-0.03^{+0.14}_{-0.15}$	$-0.06^{+0.17\pm 0.01}_{-0.18\pm 0.07}$
Dimensionless primary spin magnitude a_1	$0.32^{+0.45}_{-0.28}$	$0.31^{+0.51}_{-0.27}$	$0.31^{+0.48\pm 0.04}_{-0.28\pm 0.01}$
Dimensionless secondary spin magnitude a_2	$0.57^{+0.40}_{-0.51}$	$0.39^{+0.50}_{-0.34}$	$0.46^{+0.48\pm 0.07}_{-0.42\pm 0.01}$
Final spin a_f	$0.67^{+0.06}_{-0.08}$	$0.67^{+0.05}_{-0.05}$	$0.67^{+0.05\pm 0.00}_{-0.07\pm 0.03}$
Luminosity distance D_L/Mpc	390^{+170}_{-180}	440^{+140}_{-180}	$410^{+160\pm 20}_{-180\pm 40}$
Source redshift z	$0.083^{+0.033}_{-0.036}$	$0.093^{+0.028}_{-0.036}$	$0.088^{+0.031\pm 0.004}_{-0.038\pm 0.009}$
Upper bound on primary spin magnitude a_1	0.65	0.71	0.69 ± 0.05
Upper bound on secondary spin magnitude a_2	0.93	0.81	0.88 ± 0.10
Lower bound on mass ratio q	0.64	0.67	0.65 ± 0.03
Log Bayes factor $\ln \mathcal{B}_{\text{s/n}}$	288.7 ± 0.2	290.1 ± 0.2	—

Sensitivity evolution



An event as the GW150914 was below the threshold of $S/N=8$ to be detected by initial LIGO detectors.



The ~second Event: LVT151012

LVT: Ligo Virgo Trigger

2015, October, 12th

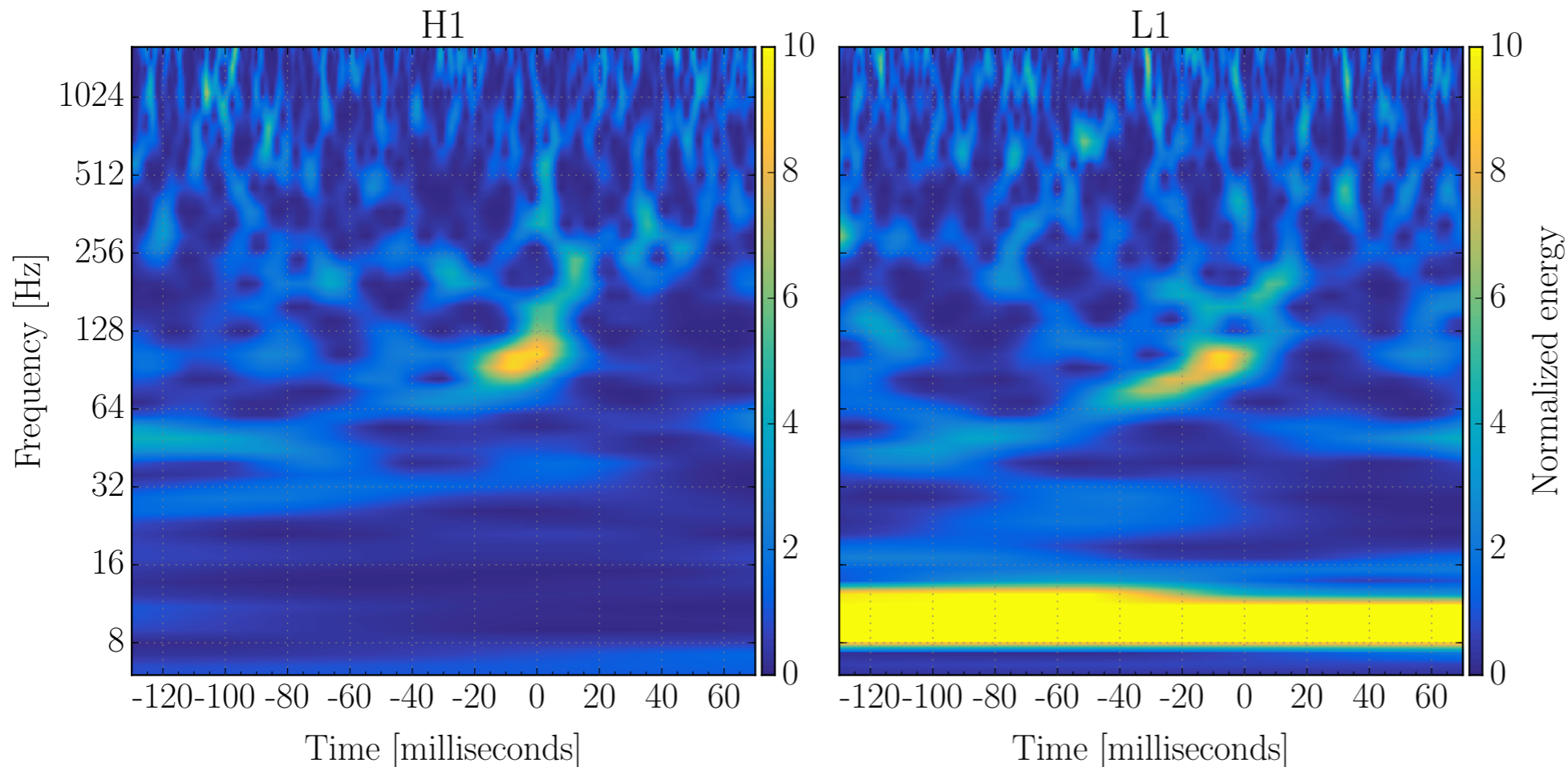
$$m_1 = 23_{-5}^{+18} M_{\odot}$$

$$m_2 = 13_{-5}^{+4} M_{\odot}$$

$$z = 0.2_{-0.1}^{+0.1}$$

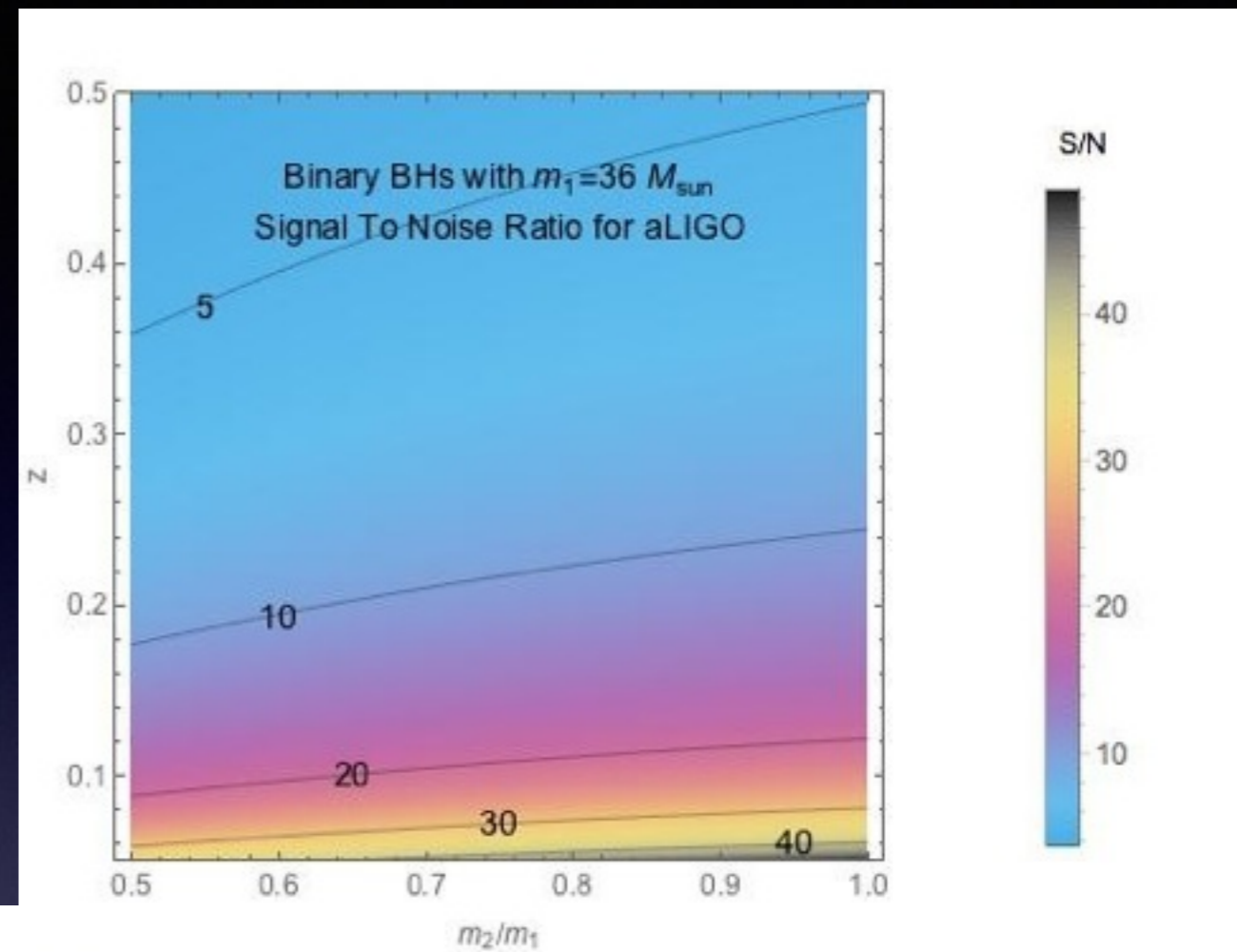
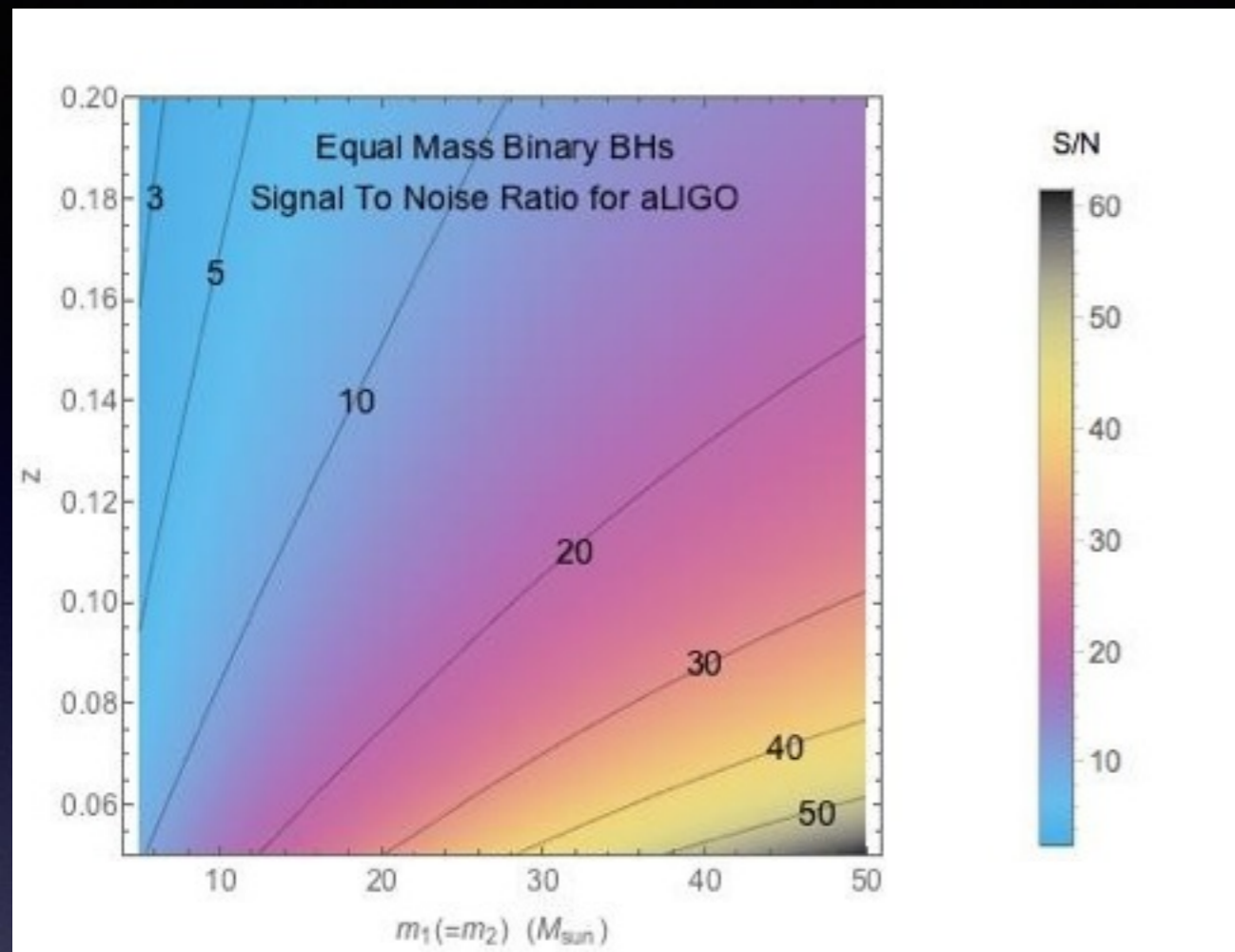
Combined S/N is 9.6 but H1
and L1 individually <8.

False rate 1 every 2.3 yrs (GW150914 was < 1 every 203000 yrs)

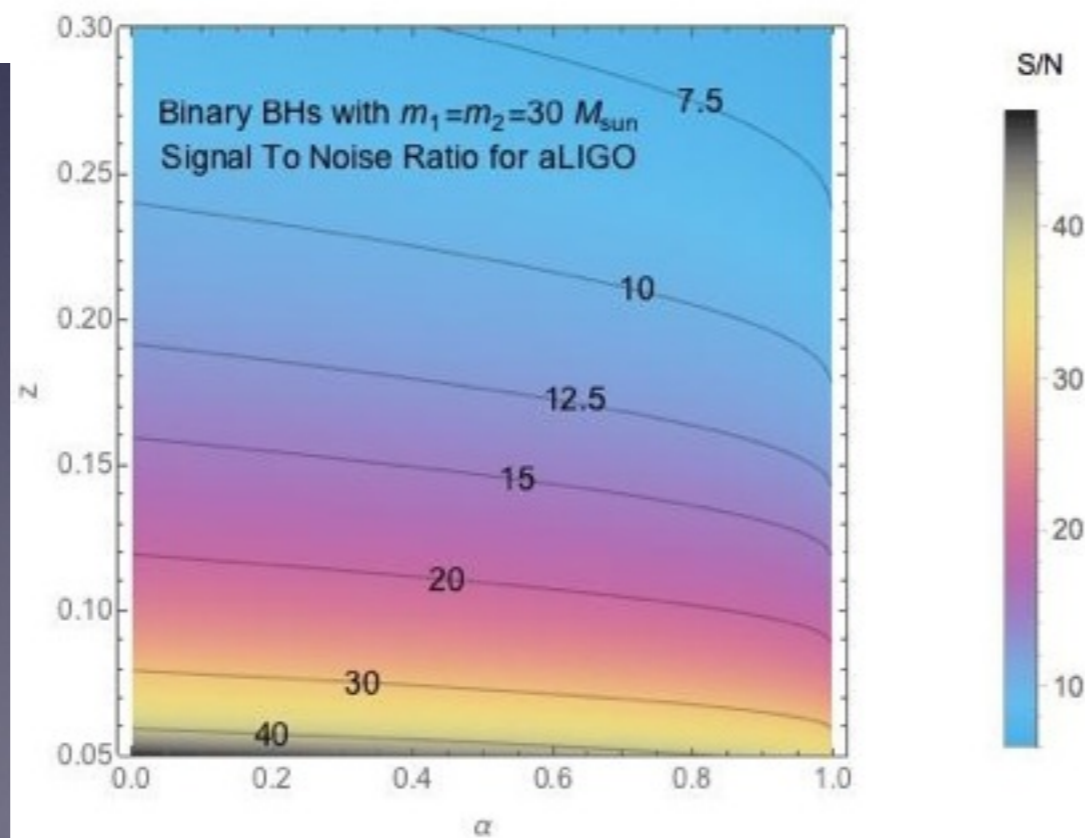


Higher Inspiral
freq. -> lower
masses

Sensitivity plots of current aLIGO in terms of sources parameters

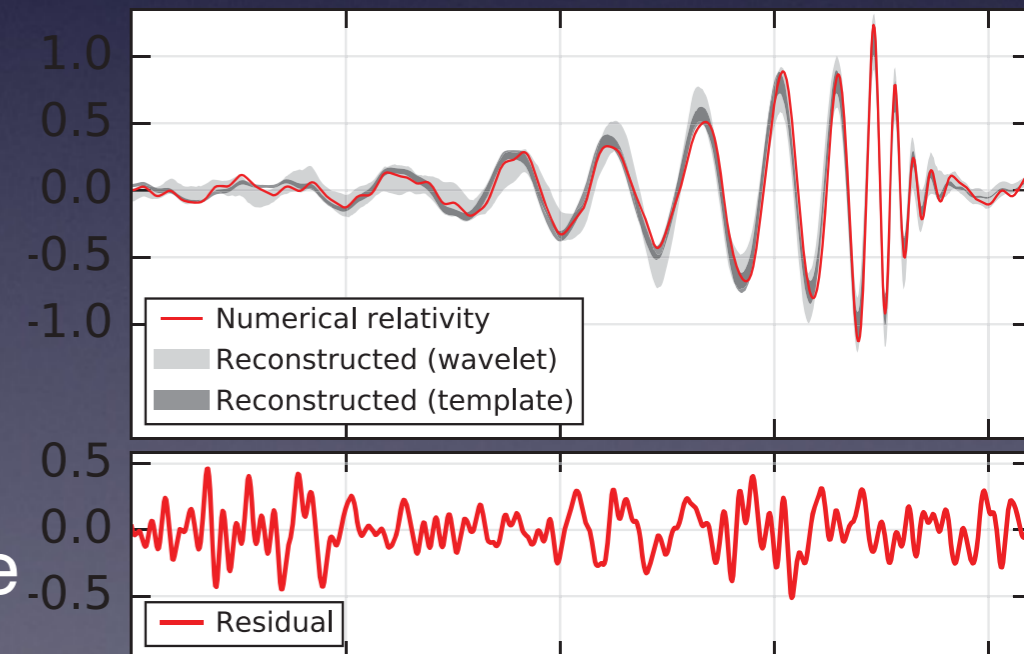


$M_c \propto \dot{f}^{-11/5}$
for fixed \dot{f}
assuming 5% of
initial mass goes
in GWs

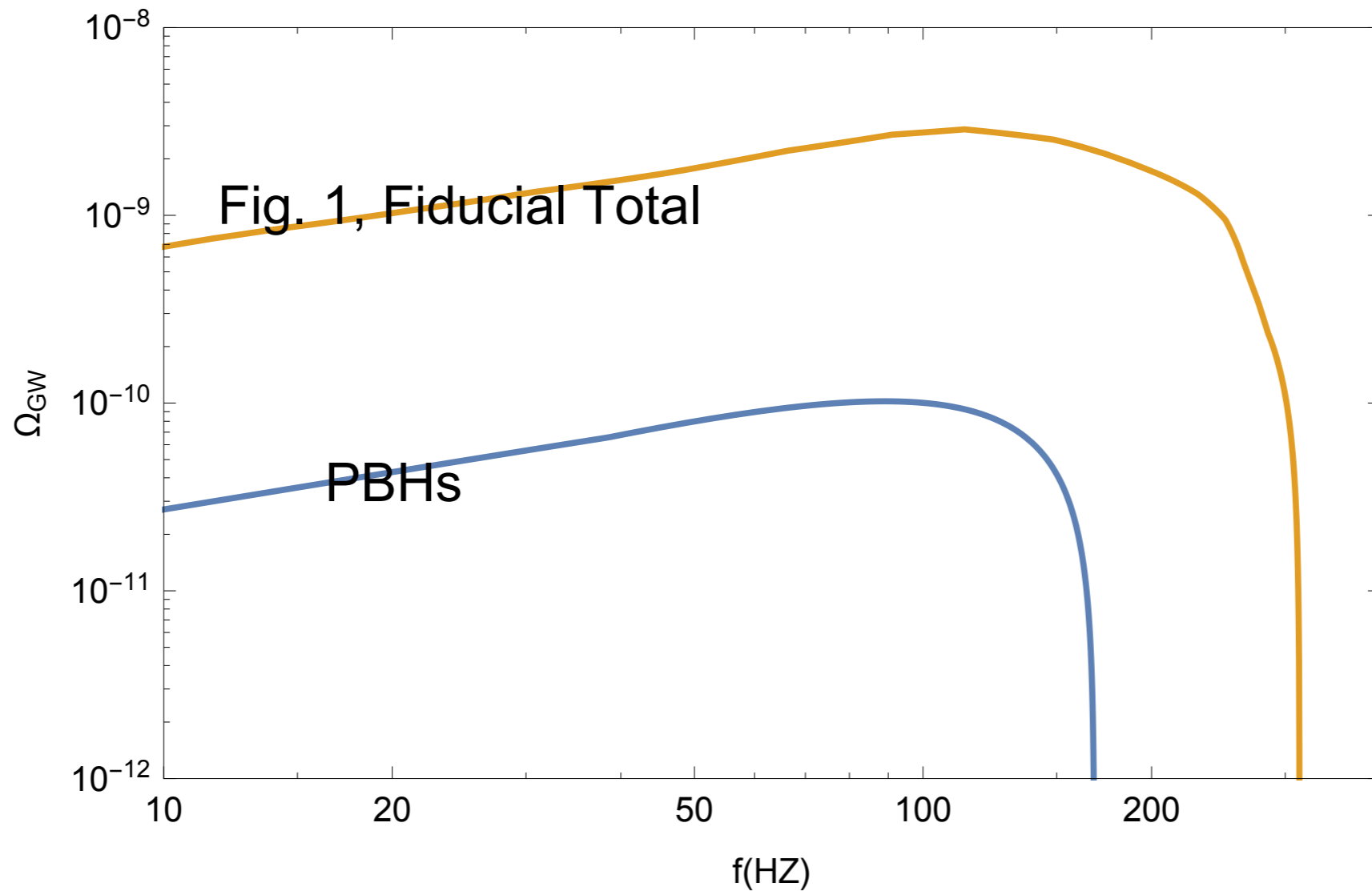


Searches for a signal

- Using waveforms (or “templates”) of merging compact objects (250000 templates), $1 - 99 M_{\odot}$ and $\alpha \in (0, 0.99)$
- searching for transient signals (using linear combinations of Sine-Gaussian wavelets). GW150914 was detected by both methods.
- LIGO measures, frequency-range during the inspiral phase
- f_{merge} from the end of the inspiral phase
- $\dot{f} \equiv df/dt$ during the inspiral phase
- h_c during the inspiral and merger phases
- $f_{ring\ down}$ from the end of the merger phase
- S/N main contribution from the merger phase but also some from the insp.



Assuming indeed LIGO will be able to study the stochastic grav. wave background then:



No power at high frequencies from PBHs